



Relaxation rate in ultra thin silicon membrane

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CREATING A DIMENSION OF INFINITE POSSIBILITIES



Motivation

Confined acoustics phonons



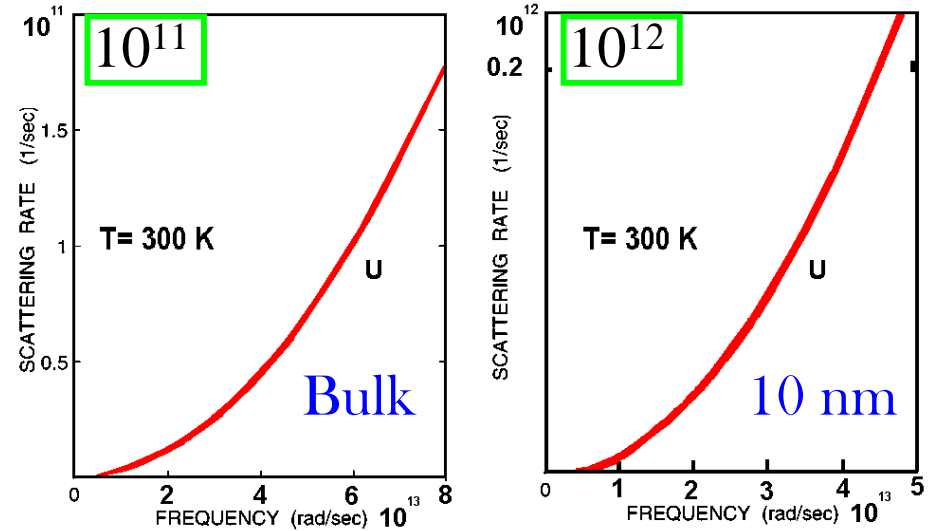
Symmetric



Antisymmetric

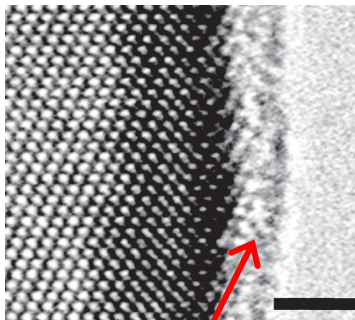


Modification of phonon lifetimes

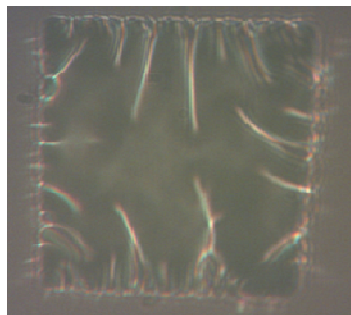


A. Balandin et al, PRB 58(1998) 1544

Effect of native oxide on dispersion relation

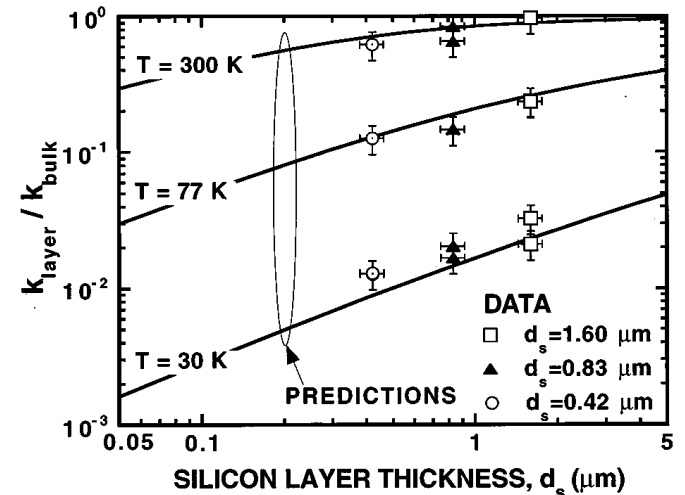


Native oxide



30 nm

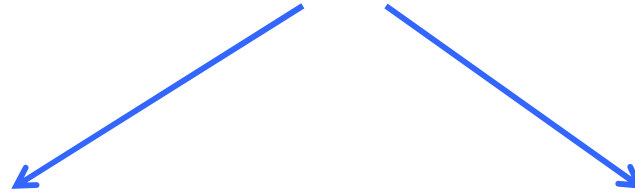
Heat transport in nm scale



M. Asheghi APL, 71(1997), 1798

Motivation

Modification of dispersion
relation (phonon engineering)



Modification of group velocity → Modification of relaxation rate



Thermal conductivity



Improve ZT

Outline

♣ Theory:

- Elastic continuum model
- Equation of motion
- Dispersion relation

♣ Relaxation rate & Thermal conductivity:

- Scattering mechanisms: Umklapp & Normal
- Relaxation rate for U-processes
- Thermal conductivity

♣ Numerical results: discussion

♣ Conclusions

Theory

Elastic continuum theory

- Stress-strain relationship
- “Macroscopic model” i.e. **continuum-based** and **isotropic**
- Well-suited for nanoscale confinement studies

Stress tensor

Elastics Tensor

Isotropic media

$$T_{ij} = C_{ijkl} S_{kl} \quad \longrightarrow \quad T_{ij} = \lambda S_{kk} \delta_{ij} + 2\mu S_{ij}$$

Strain tensor

Motion equation :

$$\partial^2 U / \partial t^2 = S_T^2 \nabla^2 U + (S_L^2 - S_T^2) \nabla (\nabla \cdot U)$$

$$S_T = \sqrt{(\lambda + 2\mu) / \rho}$$

$$S_L = \sqrt{\mu / \rho}$$

U = Vector amplitude of displacement

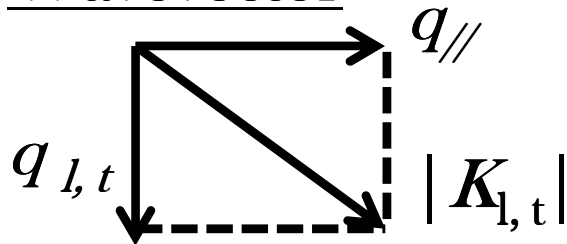
ρ = Density

λ, μ = Lamé constants

S_T, S_L = Transversal and longitudinal sound speed

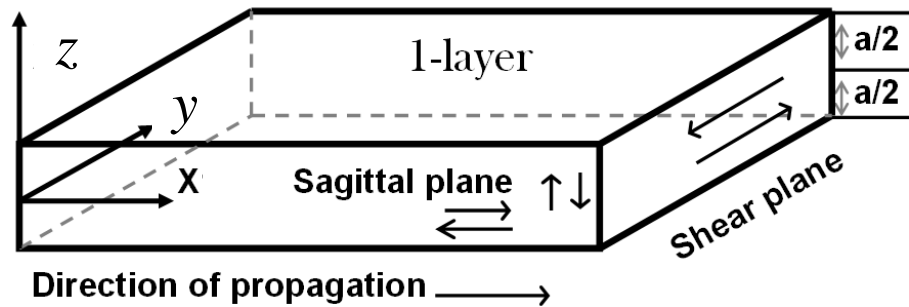
Single Layer Solutions

Wavevector



$$\partial^2 U / \partial t^2 = S_T^2 \nabla^2 U + (S_L^2 - S_T^2) \nabla (\nabla \cdot U)$$

$$\omega = S_{L,T} |K_{l,t}| = S_{L,T} \sqrt{q_{//}^2 + q_{l,t}^2}$$



$$T_{iz} \Big|_{z=\text{surface}} = 0$$

$$i = x, y, z$$

Solutions

$$U_y(z) = A \sin(q_z z) \text{ or } B \cos(q_z z) \quad \textit{Shear Waves}$$

$$U_x(z) = iAq \cos(q_l z) + iBq_t \cos(q_t z) \quad \textit{Dilatational Waves}$$

$$U_z(z) = -Aq_l \sin(q_l z) + Bq \sin(q_t z)$$

$$U_x(z) = iAq \sin(q_l z) + iBq_t \sin(q_t z) \quad \textit{Flexural Waves}$$

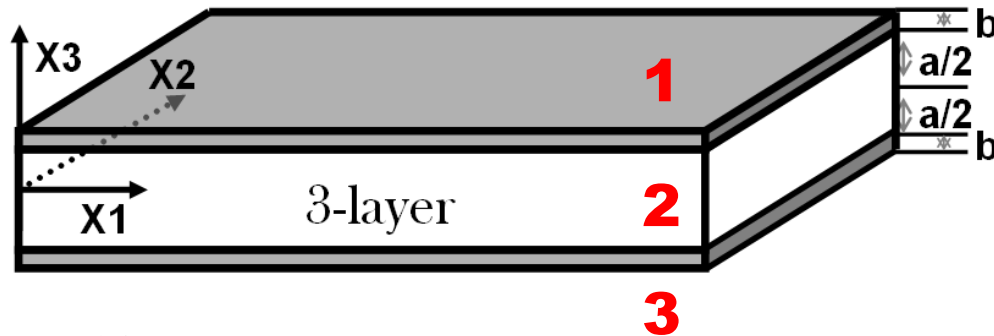
$$U_z(z) = Aq_l \cos(q_l z) - Bq \cos(q_t z)$$

Coupled solutions

Coupled solutions

Three layer Solutions

$$T_{iz} \Big|_{z= \text{surface}} = 0$$



$$T_{iz}^1 \Big|_{z= \text{interface}} = T_{iz}^2 \Big|_{z= \text{interface}}$$

$$U_i^1 \Big|_{z= \text{interface}} = U_i^2 \Big|_{z= \text{interface}}$$

$$i = x, y, z$$

$$U_x^{1,3}(z) = DW^x + FW^x$$

$$U_z^{1,3}(z) = DW^z + FW^z + B.C$$

$$U_x^2(z) = DW^x \text{ or } FW^x + C.C$$

$$U_z^2(z) = DW^z \text{ or } FW^z$$

$$+ \omega = S_{T,L}^{1,2} \sqrt{q_{t,l}^{1,2^2} + q_{//}^2} = \boxed{12 \text{ equations}}$$

System symmetry \rightarrow 6 linear equations, i.e., 6x6 matrix.

$DW^{x(z)}$: Dilatational Waves solutions en x (z) components .

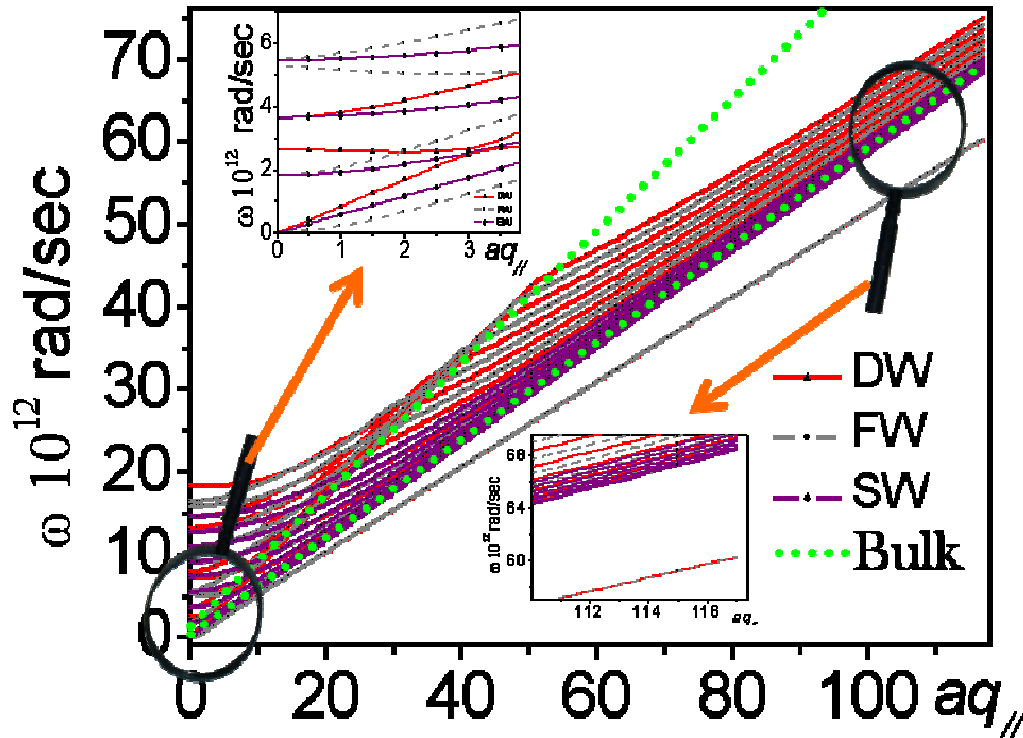
$FW^{x(z)}$: Flexural Waves solutions en x (z) components.

B. C. : Boundary Conditions.

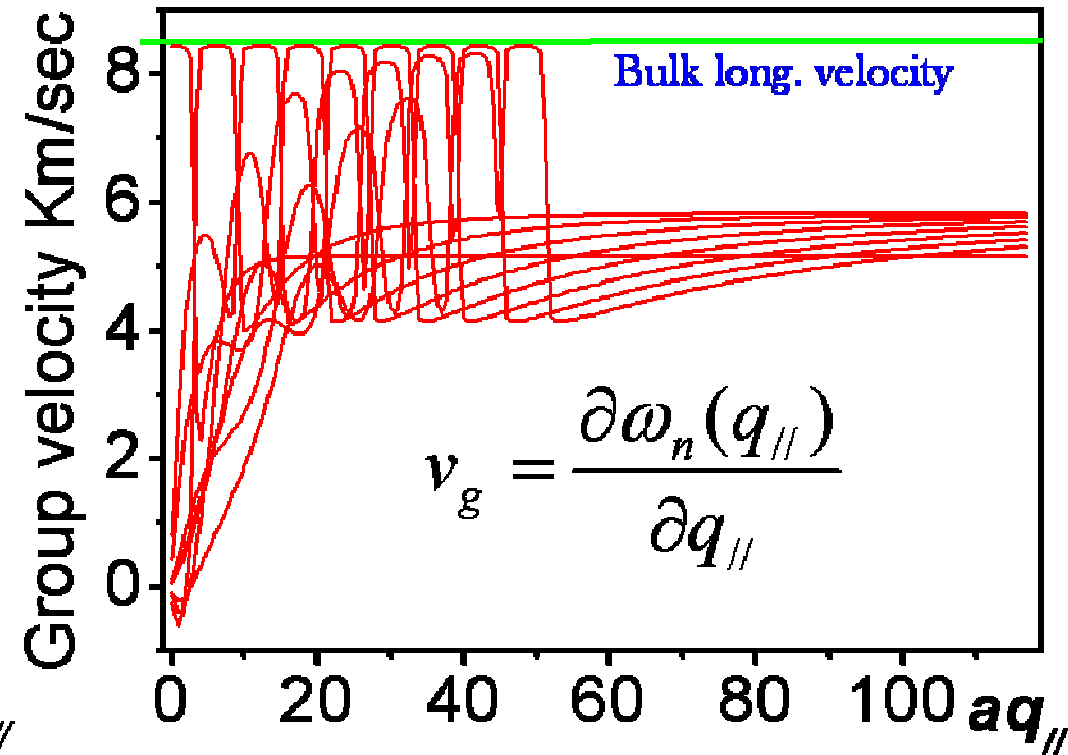
C. C. : Continuity Conditions.

Dispersion Relation & Group velocity

Sagittal & Shear Waves



Group Velocity: Dilatational Waves

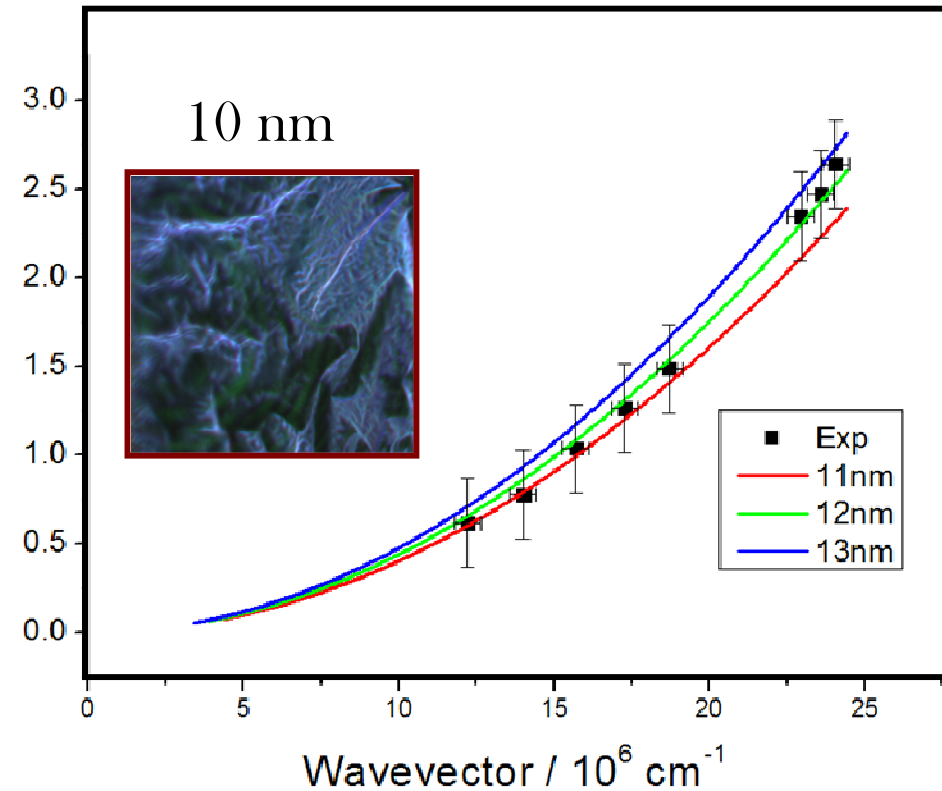
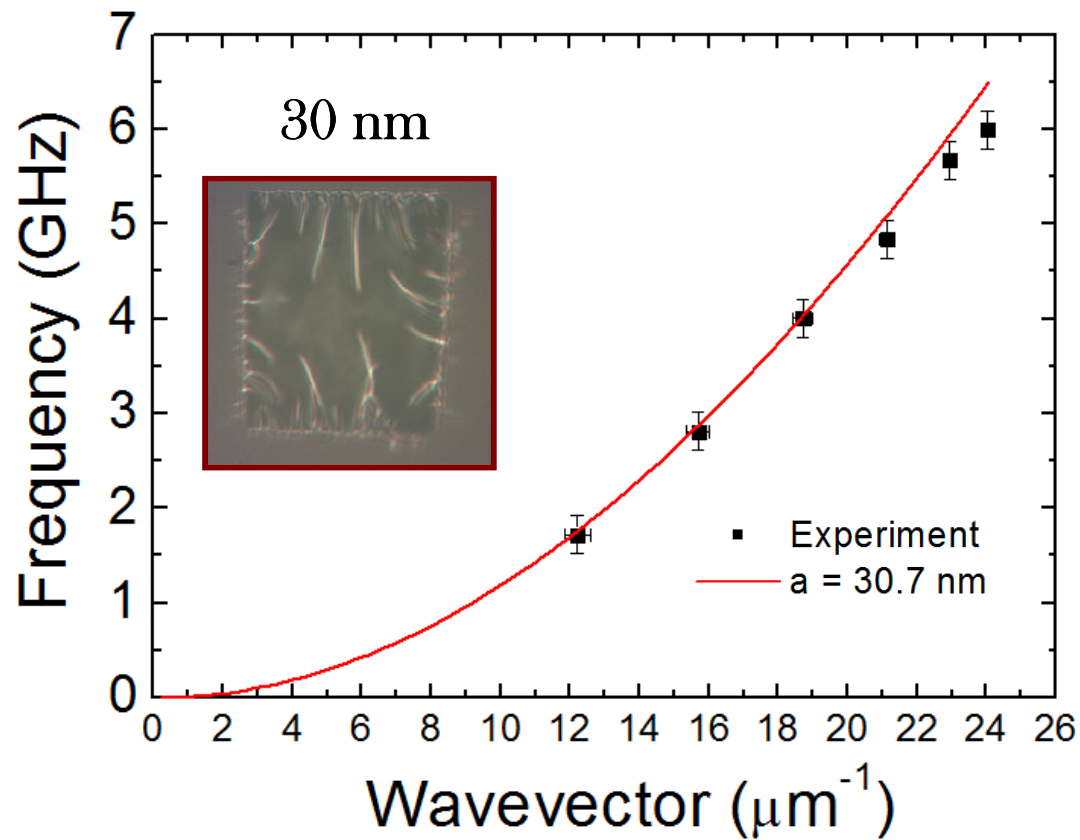


Confinement \rightarrow More modes, hybridization of sagittal waves !
 \rightarrow Group velocity (v_g) dependant on $q_{||}$

DW: Dilatational Waves
 FW: Flexural Waves
 SW: Shear Waves

Experimental Results & Theory

Flexural Waves



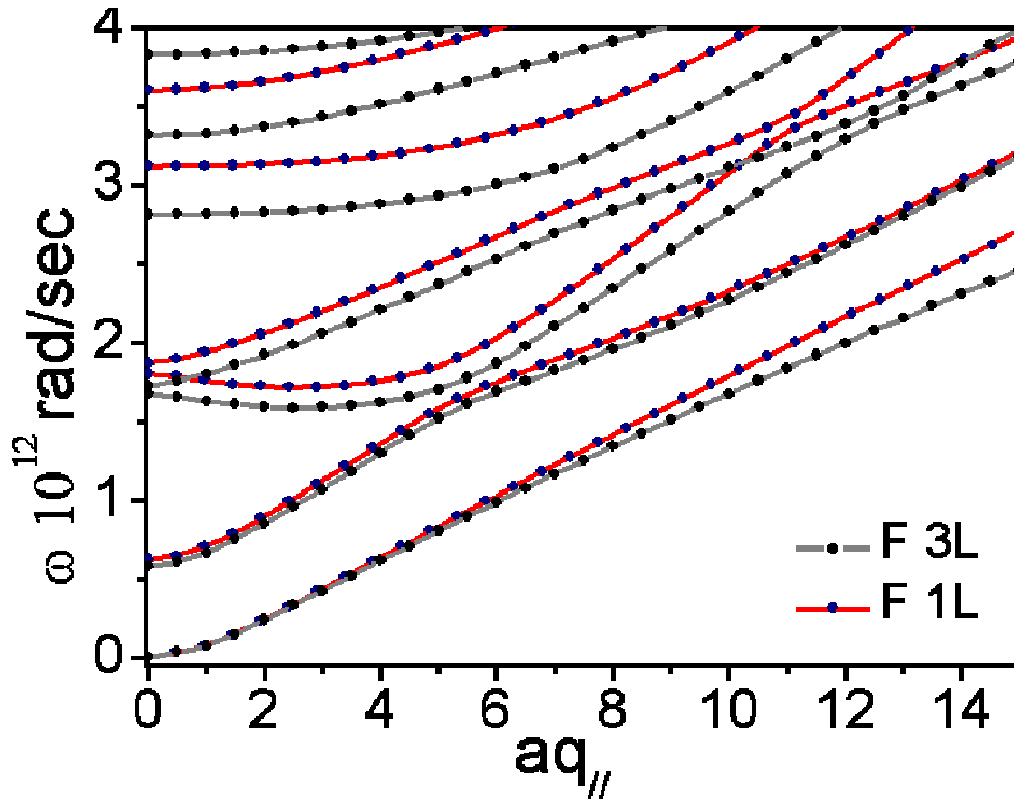
Good match between theoretical model and experiments

with

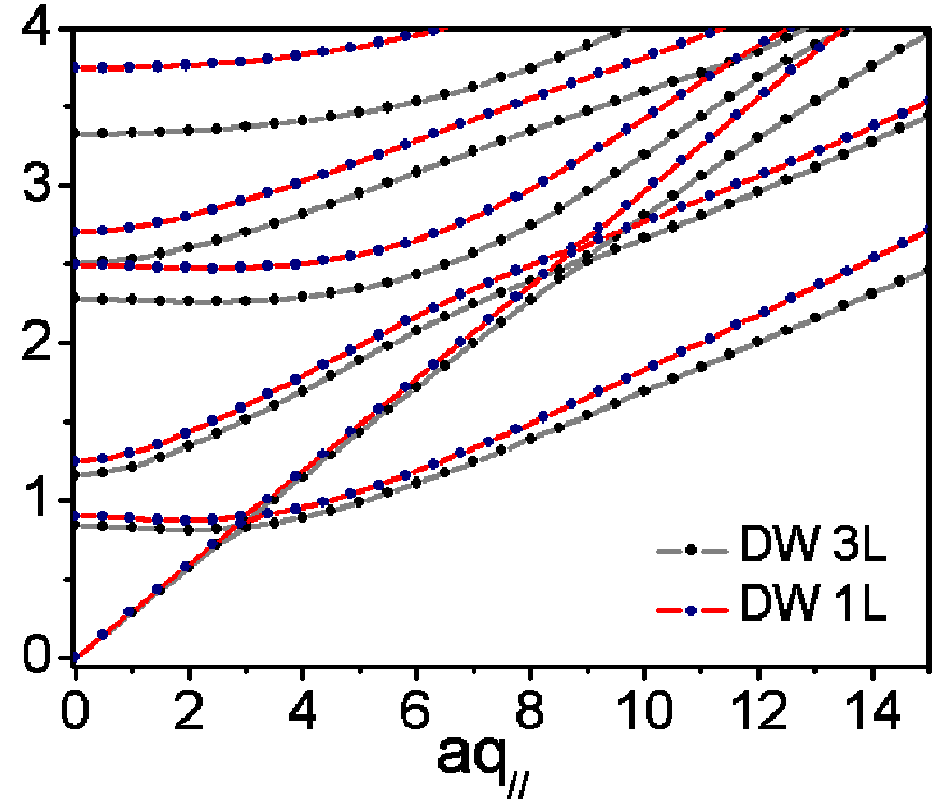


Si v/s SiO₂-Si-SiO₂

Flexural Waves



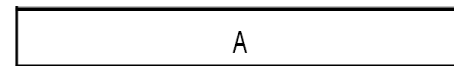
Dilatational Waves



Oxide effect: decrease of frequency for large $q_{//}$ & high order modes



$a=28.5$ nm
 $b=1.6$ nm

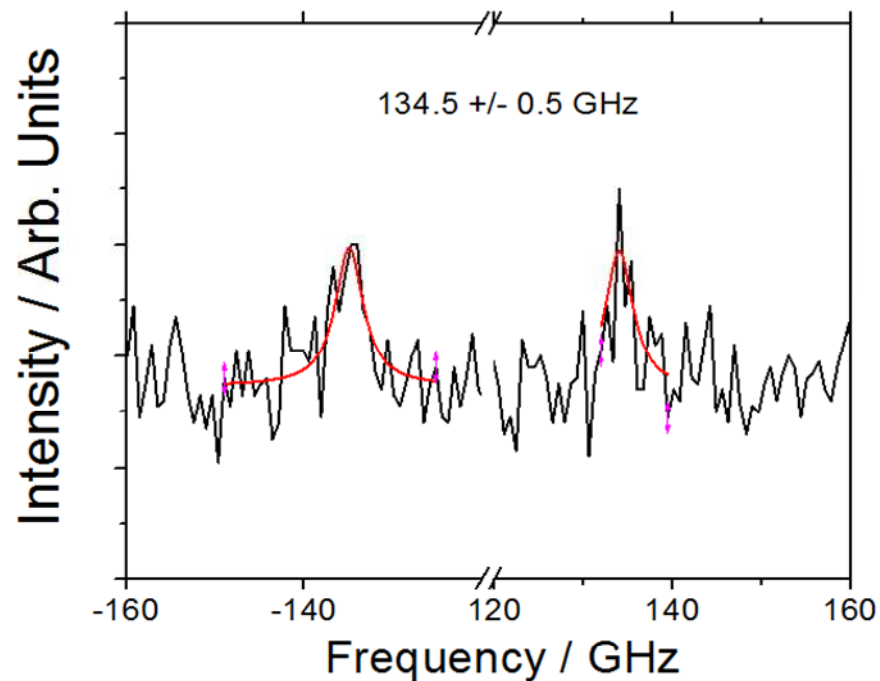


$a=29.42$ nm

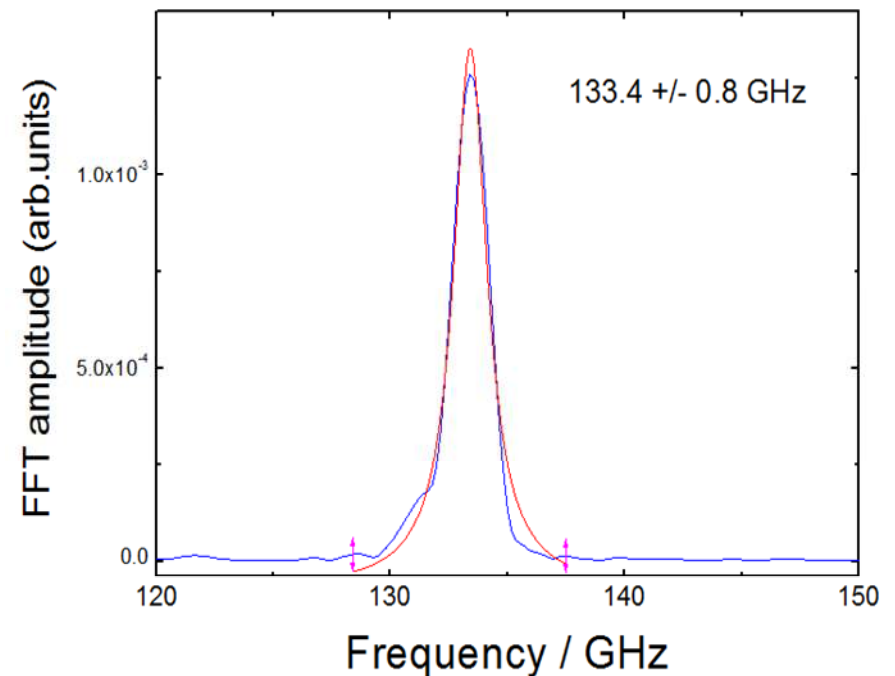
Experimental Results & Theory

30 nm Si membrane

Brillouin Scattering Spectroscopy



ASOPS technique

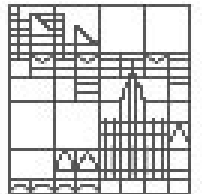


1-L Prediction: 143.5 GHz

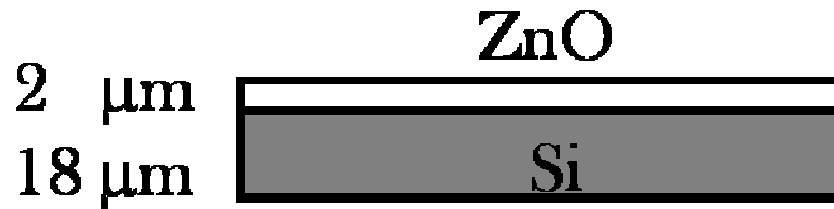
3-L Prediction: 133.8 GHz

With:

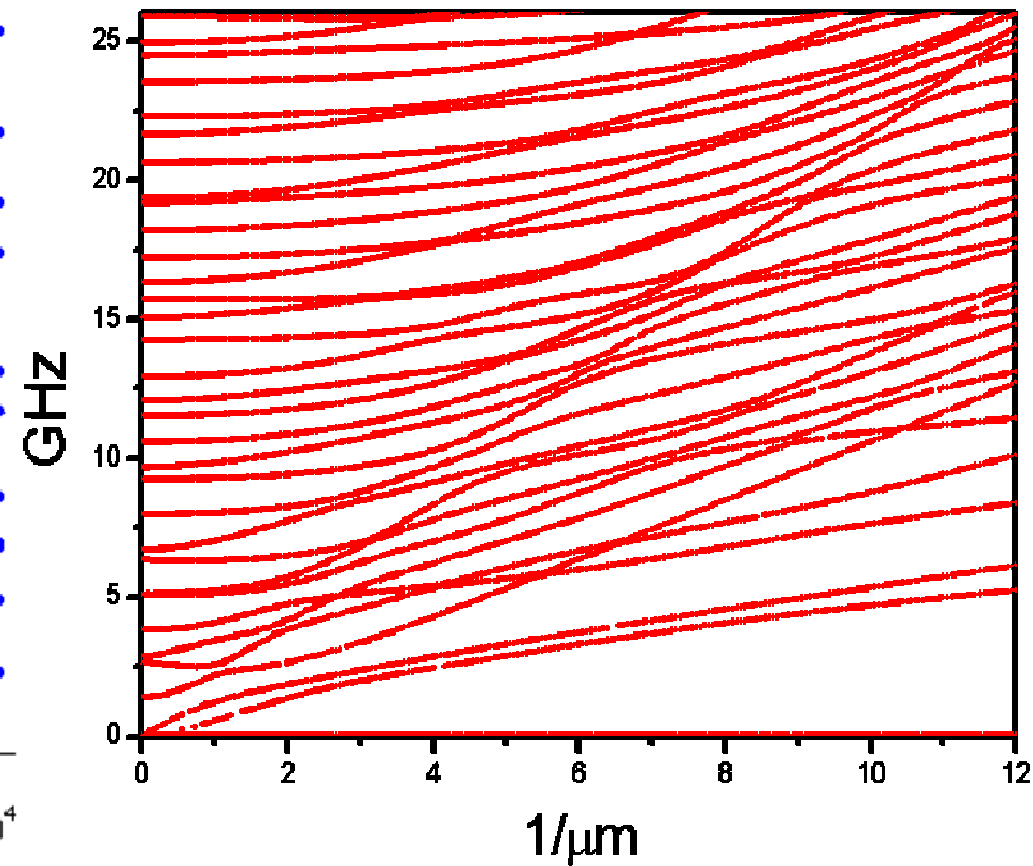
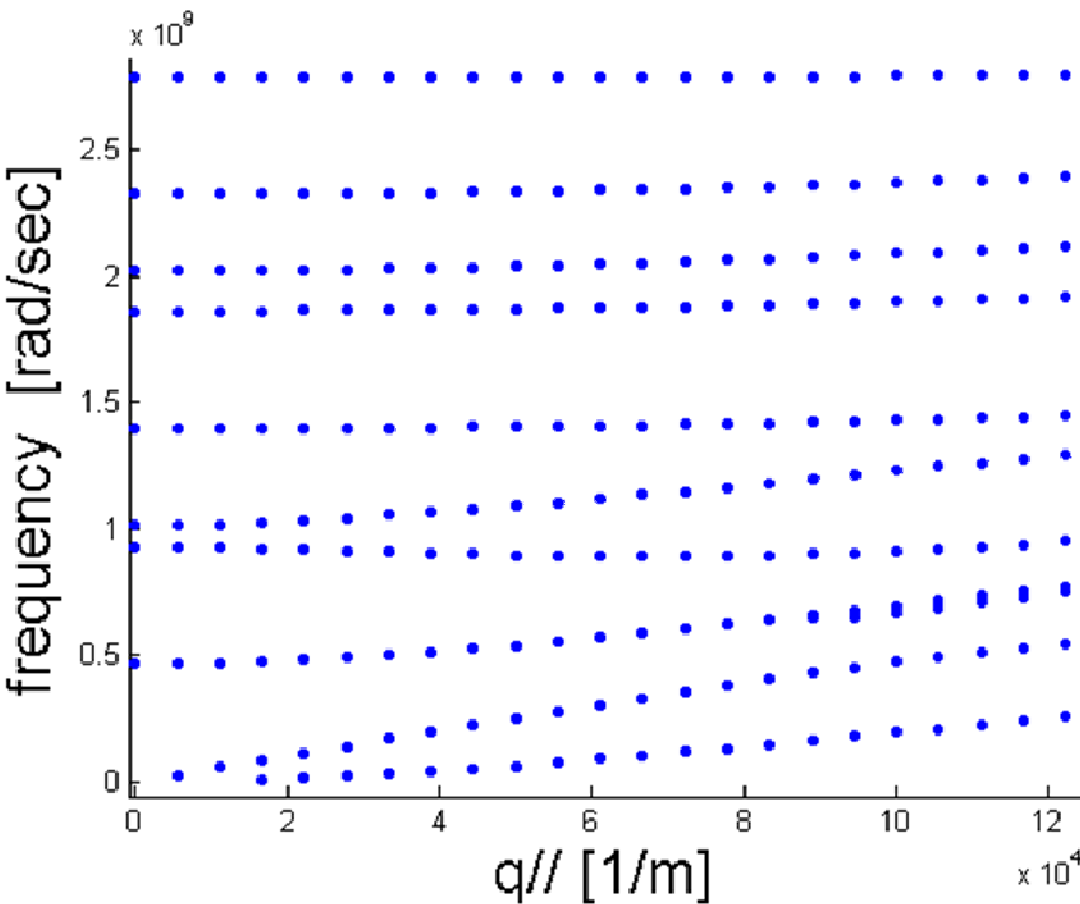
**Universität
Konstanz**



Other systems



AlN: 1,50 μm
Pt: 0,18 μm
Si: 0,40 μm



Sagittal Modes

Outline

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- Dispersion relation

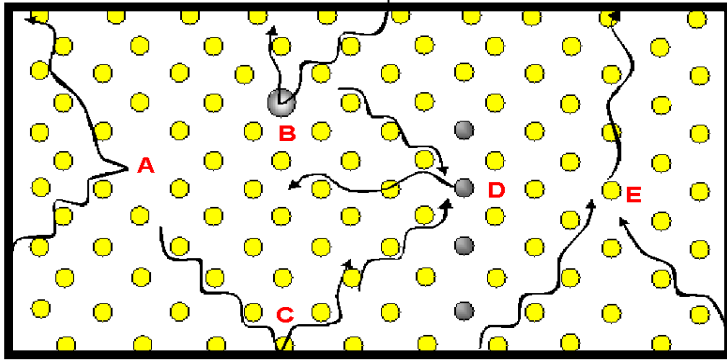
♣ Relaxation Rate & Thermal Conductivity:

- Scattering mechanisms: Umklapp & Normal
- Relaxation rate for U-processes
- Thermal Conductivity

♣ Numerical results: discussion

♣ Conclusions

Scattering Mechanisms

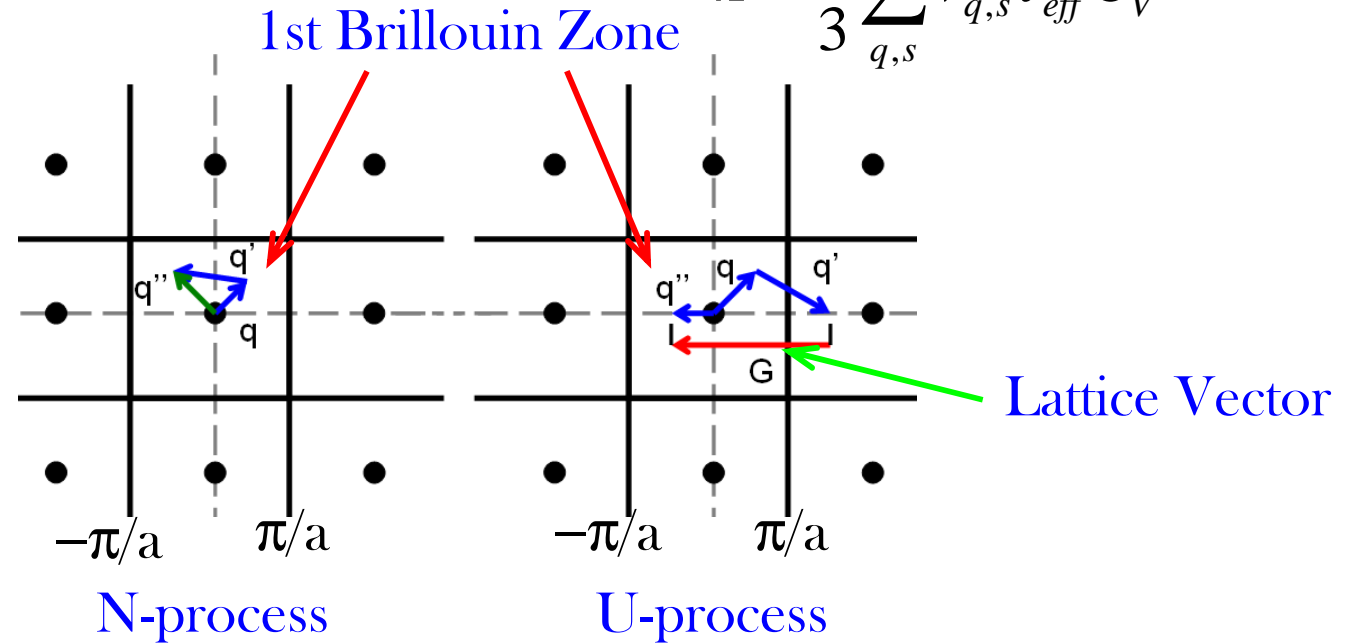
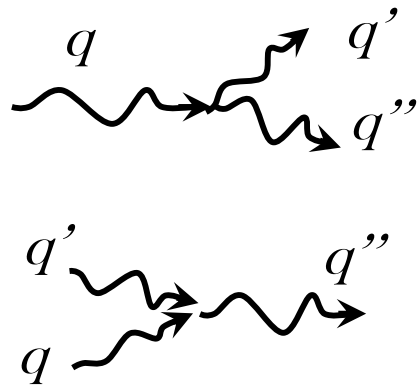


- A:** Lattice vacancy
- B:** Point defect
- C:** Crystal boundary
- D:** Umklapp
- E:** Normal

Matthiessen's rule & Thermal Conductivity

$$\frac{1}{\tau_{eff}} = \sum_j \frac{1}{\tau_j} = \frac{1}{\tau_U} + \frac{1}{\tau_b} + \frac{1}{\tau_i} + \dots$$

$$\kappa = \frac{1}{3} \sum_{q,s} v_{q,s}^2 \tau_{eff} C_V$$

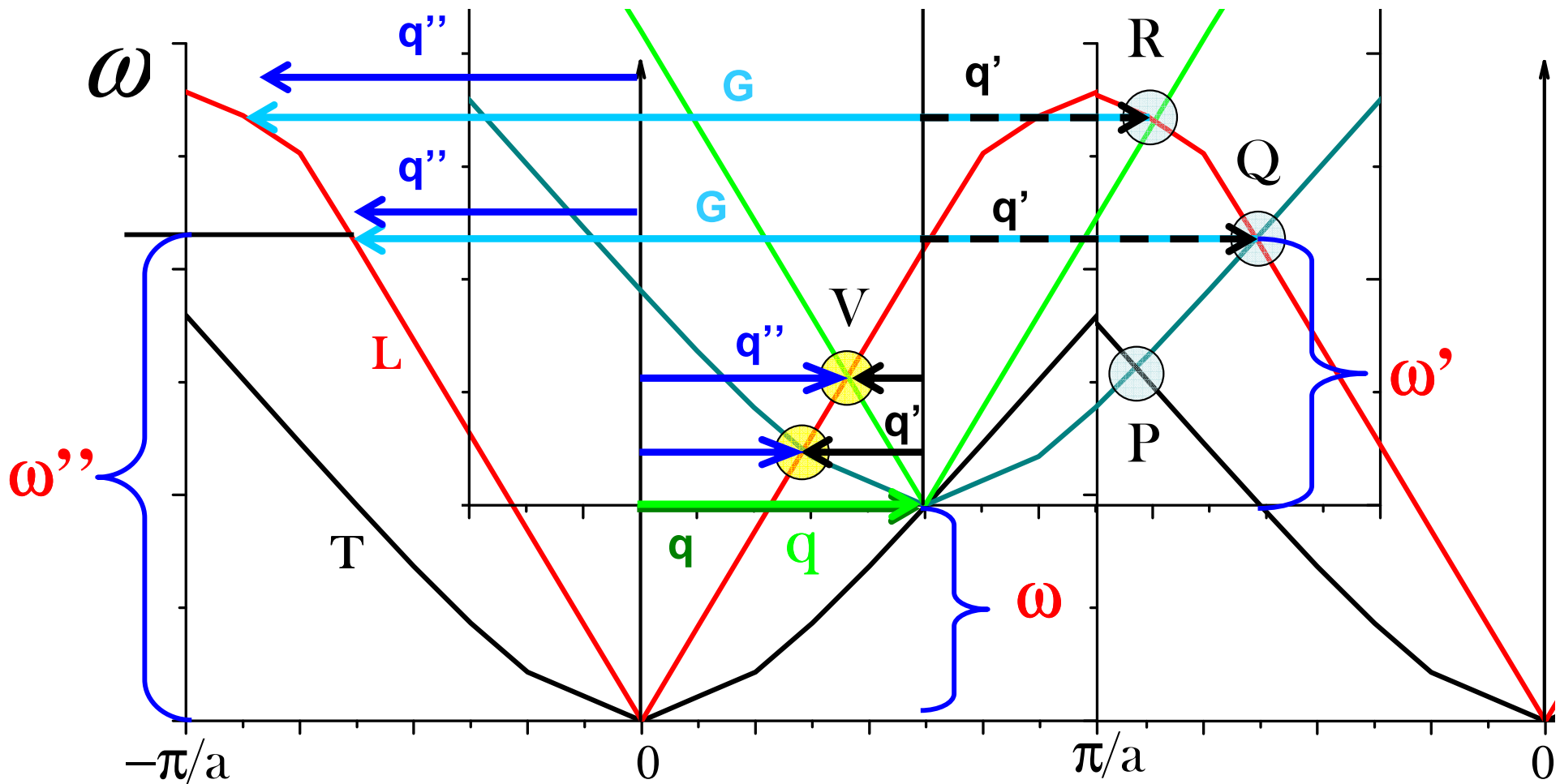


Consequence of anharmonicity → Phonon-phonon interaction → Finite phonon lifetime

Possibilities:

1. Annihilation of two phonons and creation of a third phonon.
2. Annihilation of one phonon and creation of two phonons

Selection Rules: U & N-Process



U-process $\vec{q} + \vec{q}' = \vec{G} + \vec{q}''$	$\omega + \omega' = \omega''$	N-process $\vec{q} + \vec{q}' = \vec{q}''$
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$P: T_1 + T_2 \rightarrow T_2$ If T is degenerate

$Q: T_1 + T_2 \rightarrow L$ If T is degenerate

$R: T + L \rightarrow L$

$S: T_1 + T_2 \rightarrow L$ If T is degenerate

$V: T + L \rightarrow L$

Relaxation Rate & Thermal conductivity

$$\frac{1}{\tau_{U;q,s}} = \frac{\pi\gamma^2\hbar}{\rho N_0 \Omega \bar{v}^2} \sum_{q's',q''s'',G} \omega_{q,s} \omega_{q',s'} \omega_{q'',s''} (n_{q's'} - n_{q''s''}) \delta(\omega_{qs} + \omega_{q's'} - \omega_{q''s''}) \delta_{q+q',q''+G}$$

Phonon average group velocity

$$\bar{v}_s = \frac{\sum_q v_{qs} n_{qs}}{\sum_q n_{qs}}$$

Kinetic theory

$$\kappa = \frac{1}{3} \sum_{q,s} v_{q,s}^2 \tau_{eff} C_V$$

$$\kappa = \frac{\hbar \rho \bar{v}^2}{3\pi\gamma^2 k_B T^2} \sum_{qs} v_{qs}^2 \omega_{qs}^2 n_{qs} (n_{qs} + 1) \left\{ \sum_{q's';q''s'';G} (n_{q's'} - n_{q''s''}) \omega_{qs} \omega_{q's'} \omega_{q''s''} \delta(\Delta\omega) \delta_{q+q'-q'',G} \right\}^{-1}$$

More Branches → More interactions → Decrease thermal conductivity.

Change of group velocity (v_{qs}) → Change in thermal conductivity

A. AlShaikhi and G.P. Srivastava, *Phys. Rev. B*, 76 (2007), pp. 195205-1.

B. G. P. Srivastava, “*The Physics of Phonons*”, Taylor & Francis Group, NY, 1990.

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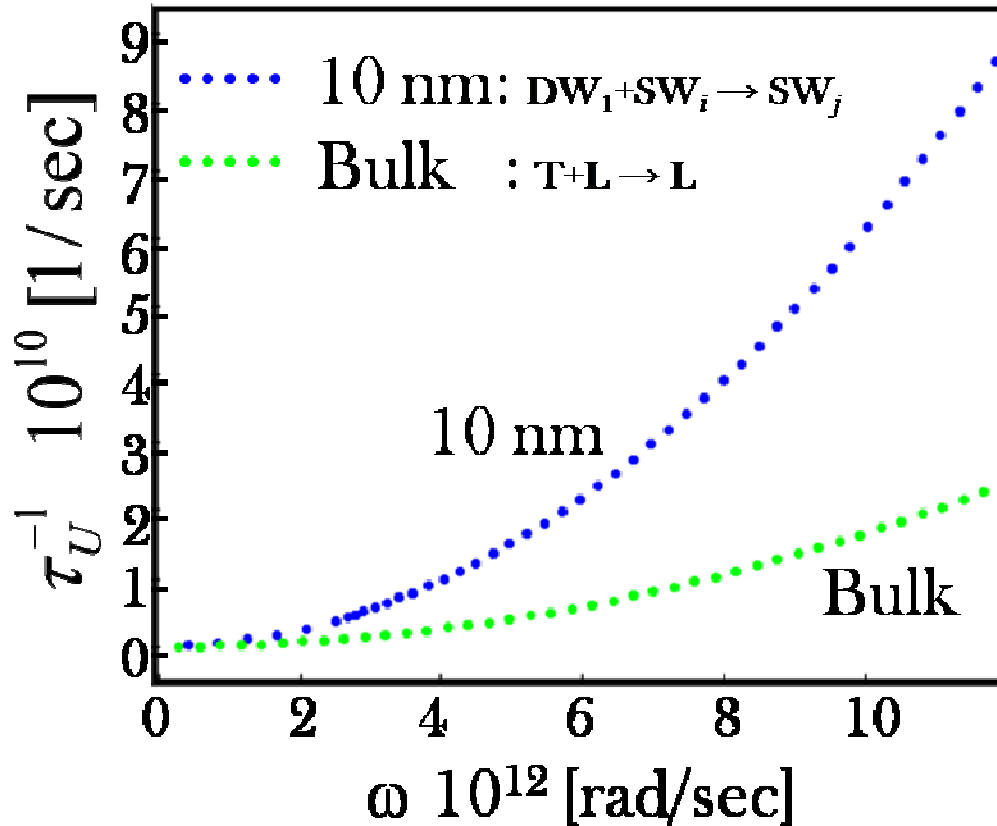
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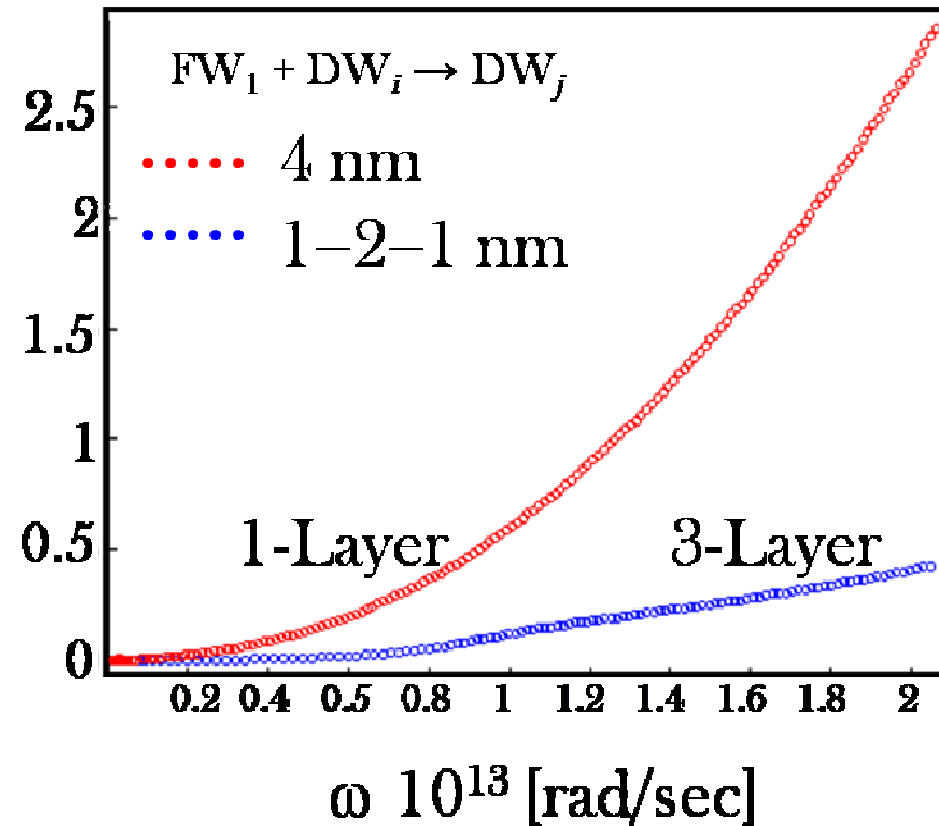
Results: Relaxation Rates

T= 300 K

Bulk & 1-Layer membrane



1 & 3 Layer membranes

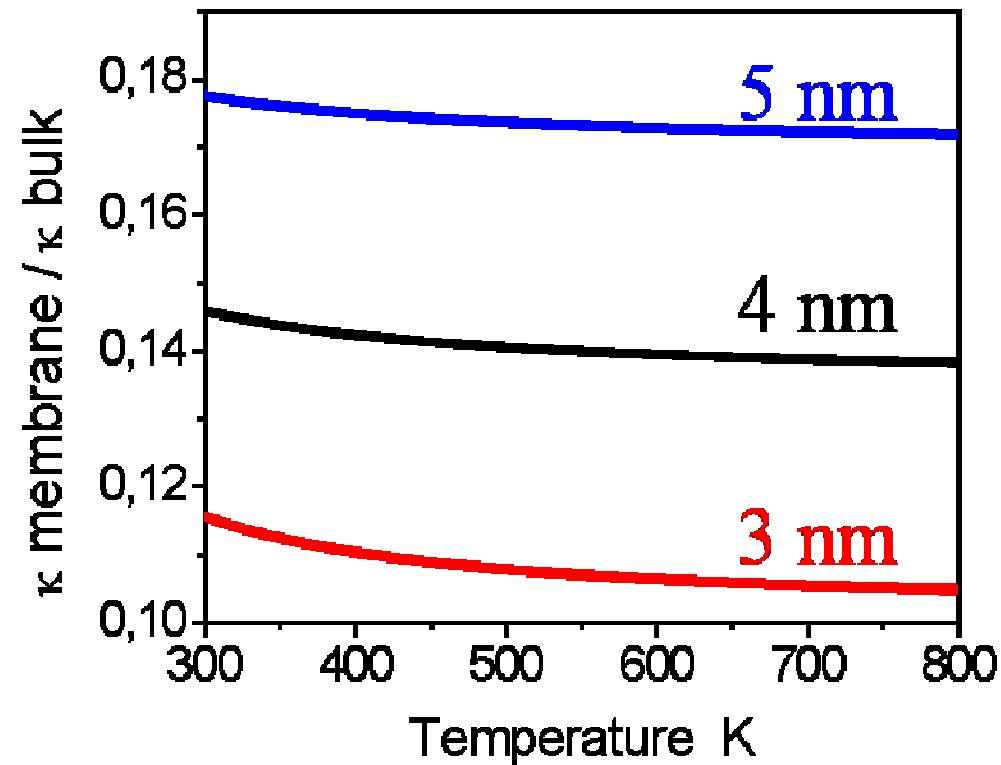
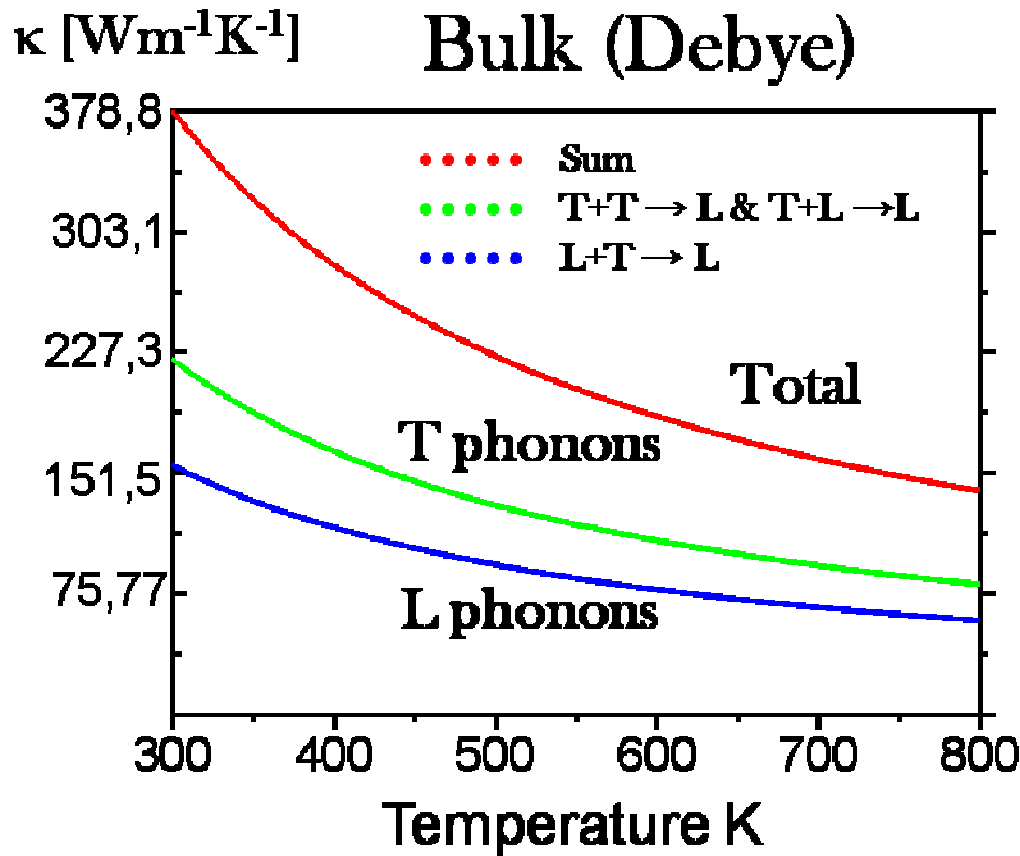


The modification of dispersion and group velocity lead to increase of relaxation rate compared with the bulk (left).

The decrease in frequency of 3-layer membrane lead to decrease on relaxation rate compared with 1-layer membrane (right).

Results: Thermal Conductivity

Intrinsic thermal conductivity



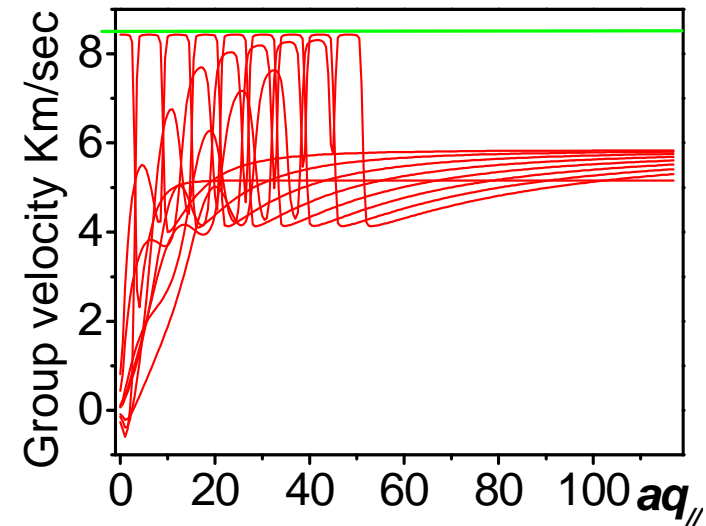
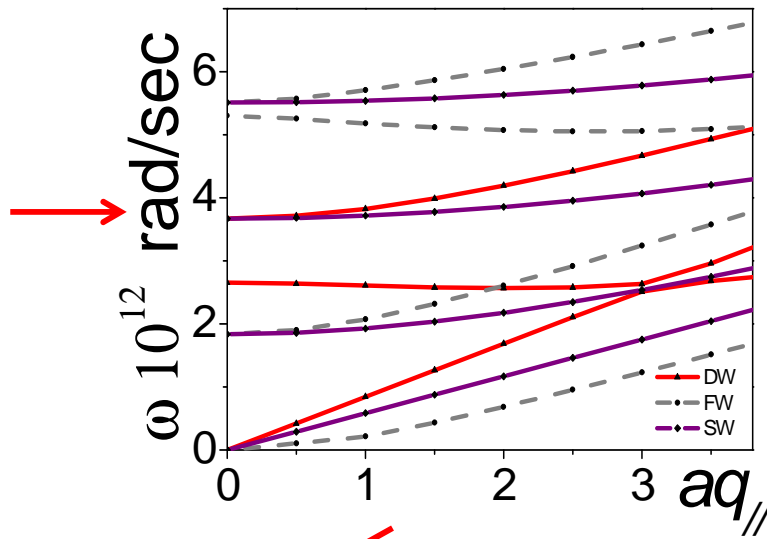
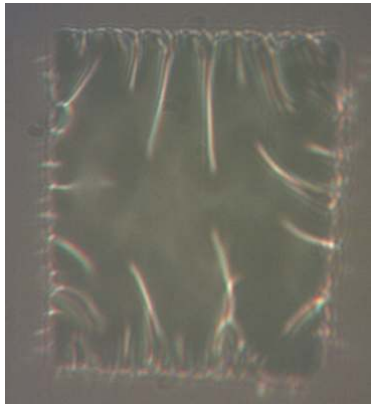
The numerical calculations predict a strong decrease of the thermal conductivity of a 1-layer membranes.

Summary

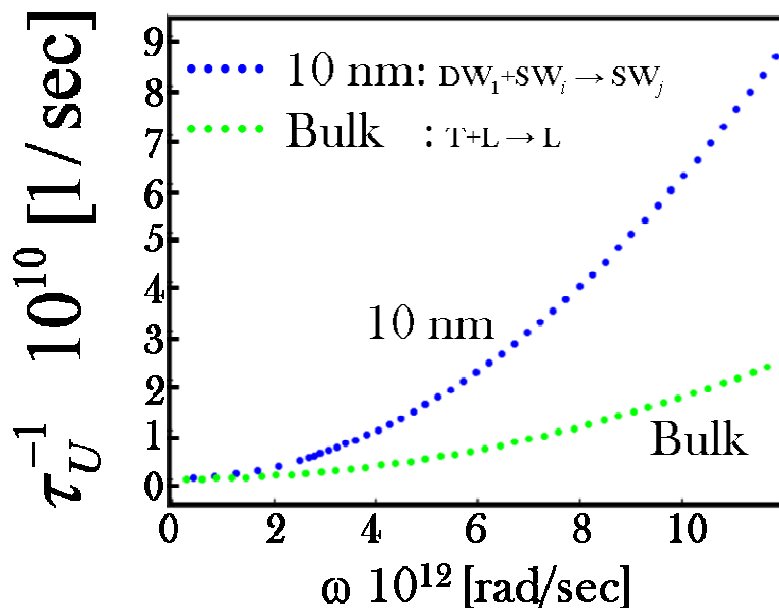
Spatial confinement

Modification of dispersion relation

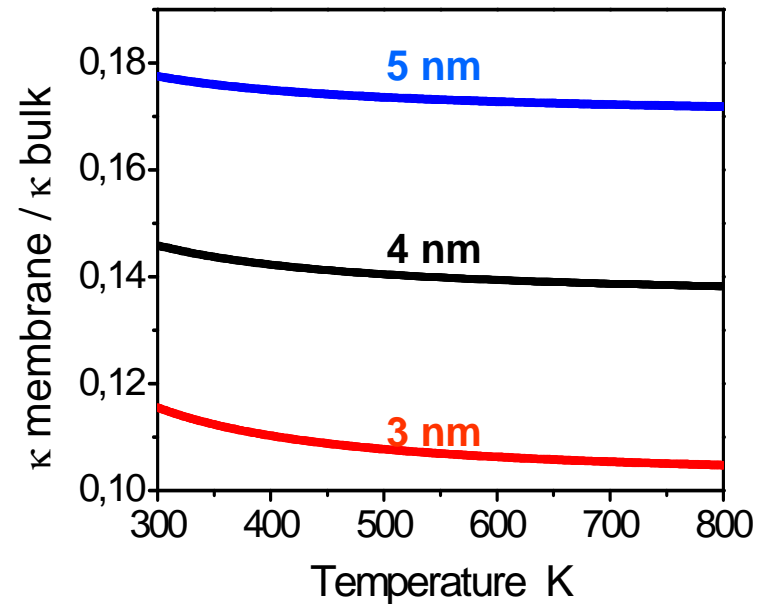
Modification of group velocity



Increase of relaxation rate



Decrease of thermal conductivity



Conclusions

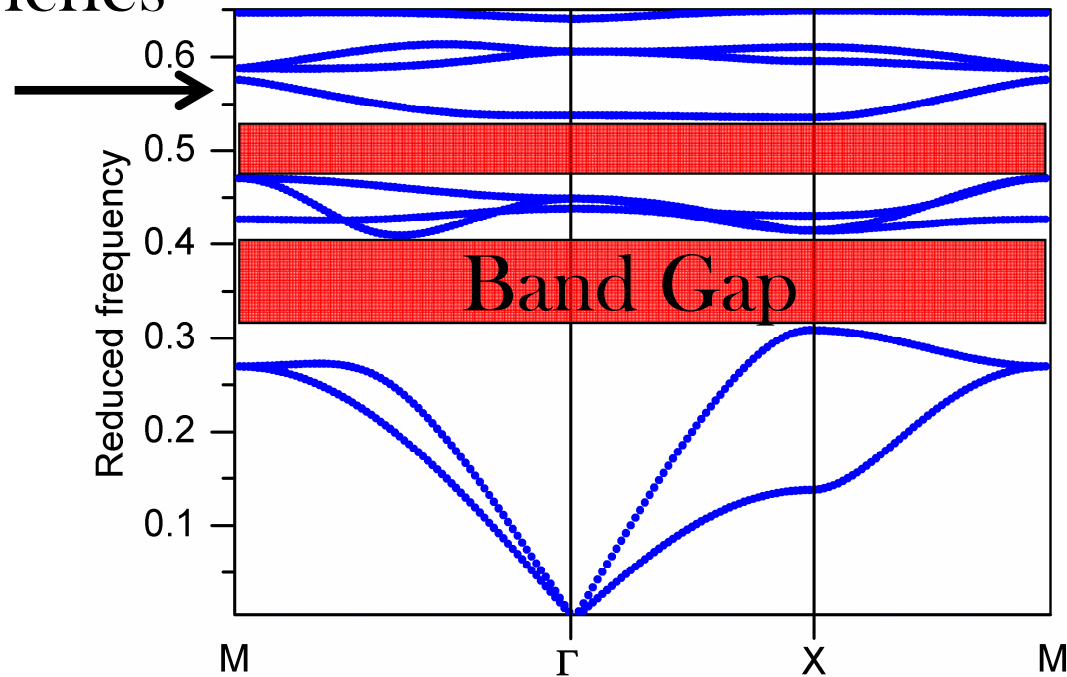
- ♣ Hybridization of sagittal waves.
- ♣ Group velocity dependence on $q_{//}$: $v_g(q_{//})$.
- ♣ Oxide effect: decrease of frequency for large $q_{//}$
→ Decrease in relaxation rate.
- ♣ The change in dispersion relation and the emergence of more branches implies an increase of phonon-phonon interactions. This causes an increase in relaxation rates and a corresponding decrease in the thermal conductivity.

Future Work

Phononic structures

PWE solutions

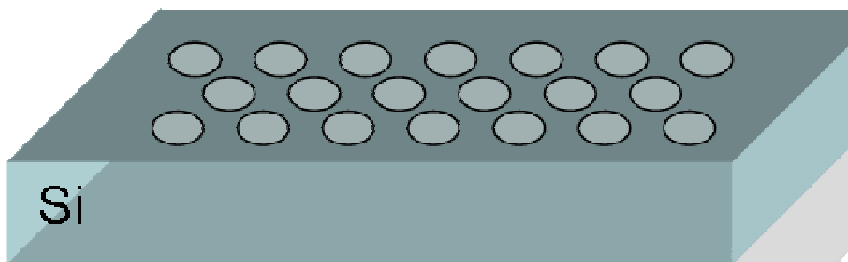
Flat Branches



With:

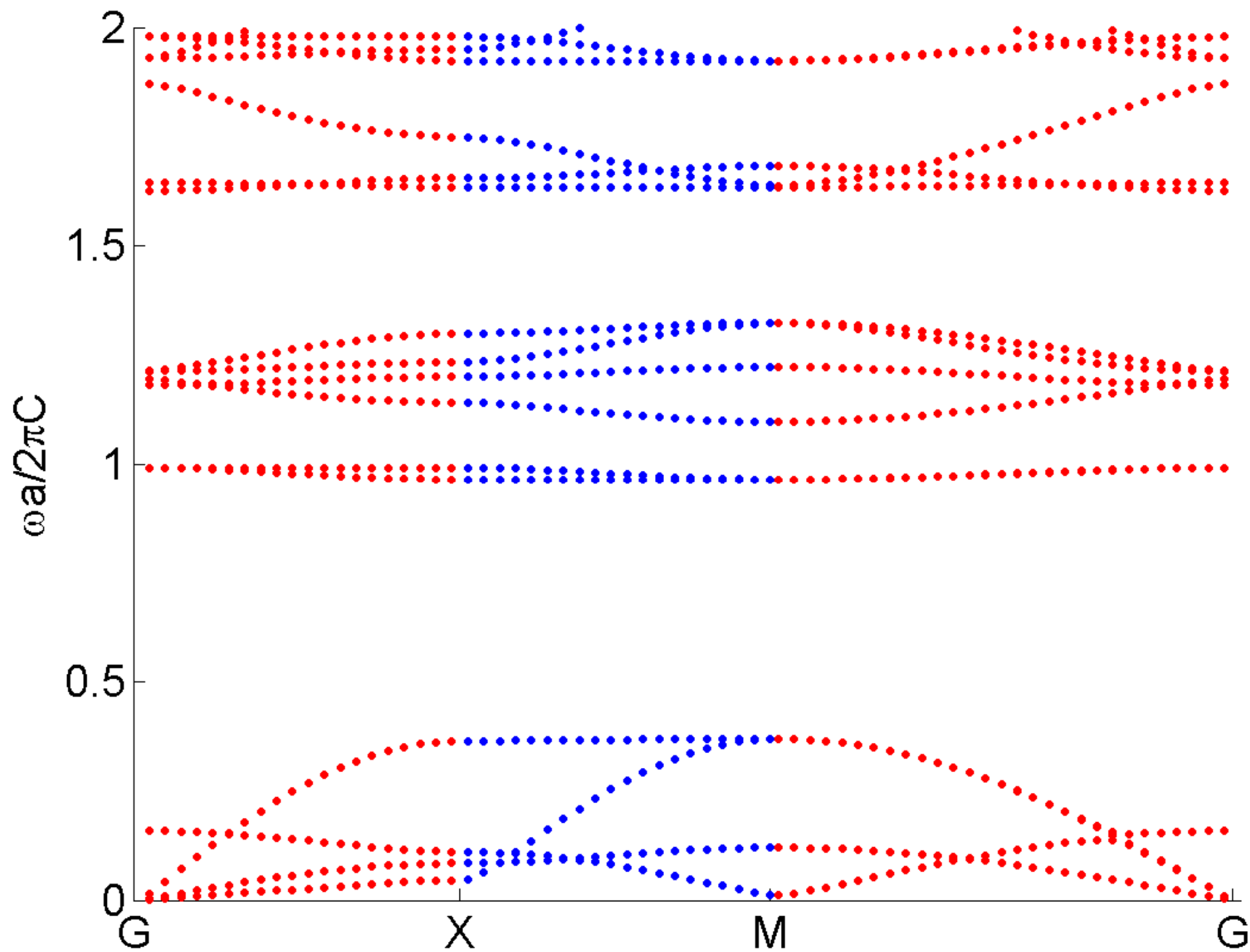


Theoretical 2D dispersion relation for carbon cylinder in an epoxy matrix (sagittal waves)



Y. Pennec et al, surface science reports
65(2010), pp 229-291

R/A = 0.5000 ff = 0.7854



Thanks you