

Elogio del rumore

dalla

Risonanza Stocastica

al

Energy Harvesting

Luca Gammaitoni

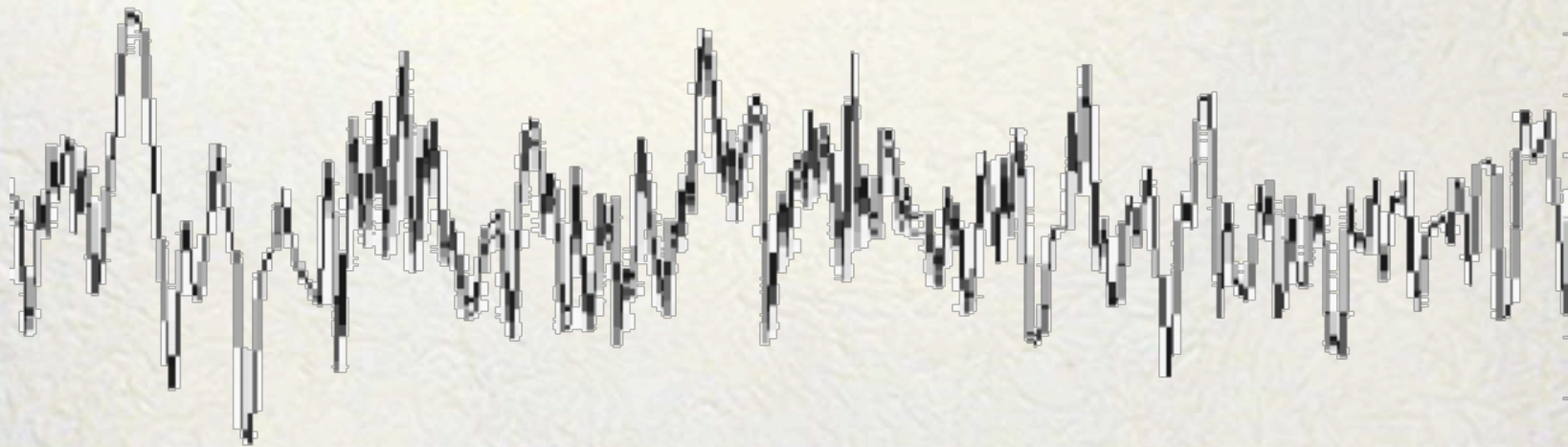
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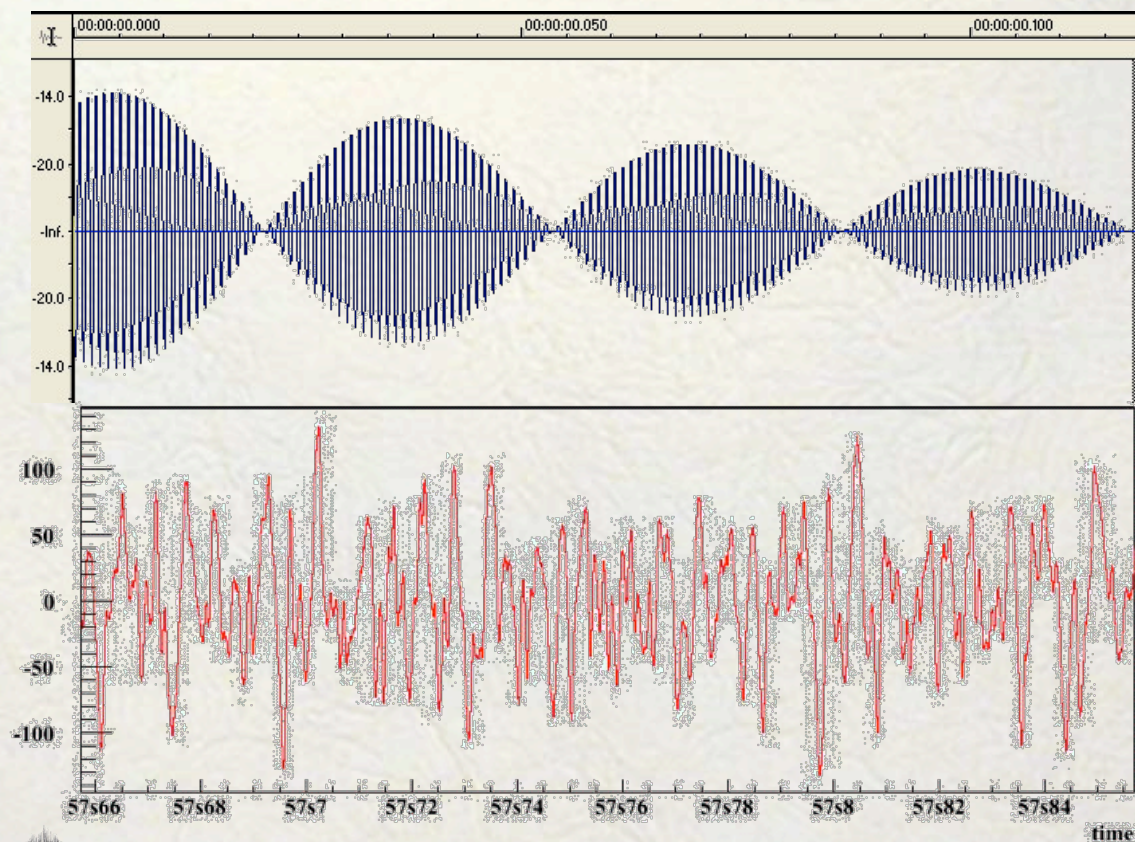
*La metta così:
dove niente sta al posto giusto, c'è disordine.
Dove al posto giusto non c'è niente, lì c'è ordine*

B. Brecht, Dialoghi di profughi



*La metta così:
dove niente sta al posto giusto, c'è disordine.
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B. Brecht, Dialoghi di profughi



What do we mean with **noise** ?

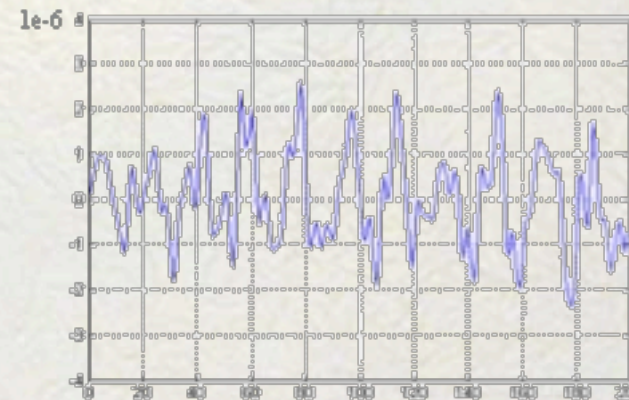
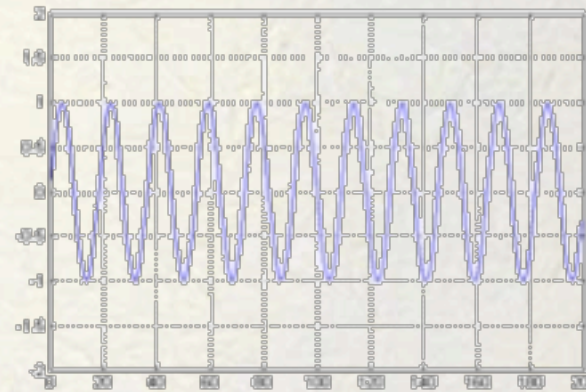
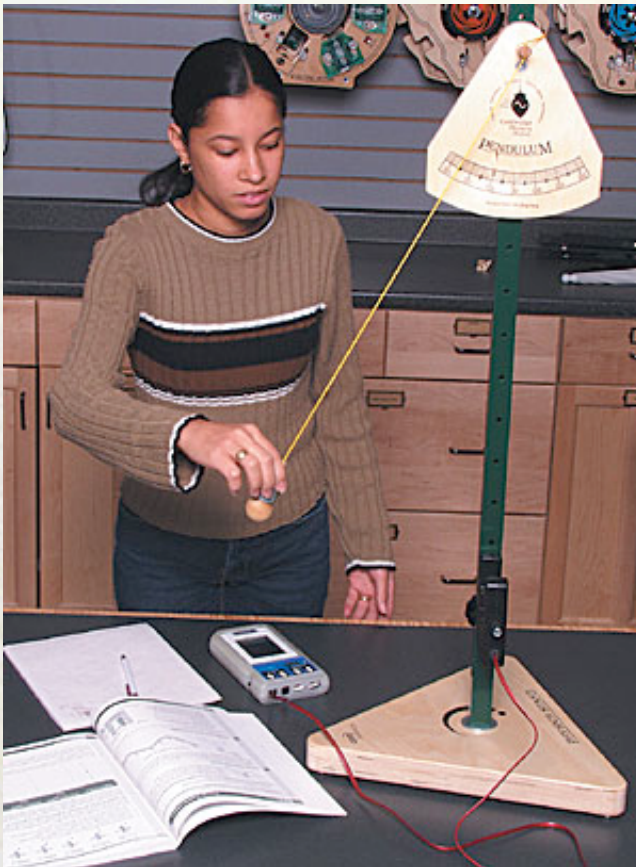
Noise in Physics means:

A disturbance, especially a random and persistent disturbance, that obscures or reduces the clarity of a signal.

Noise becomes relevant when we try to measure a physical quantity with high **precision**

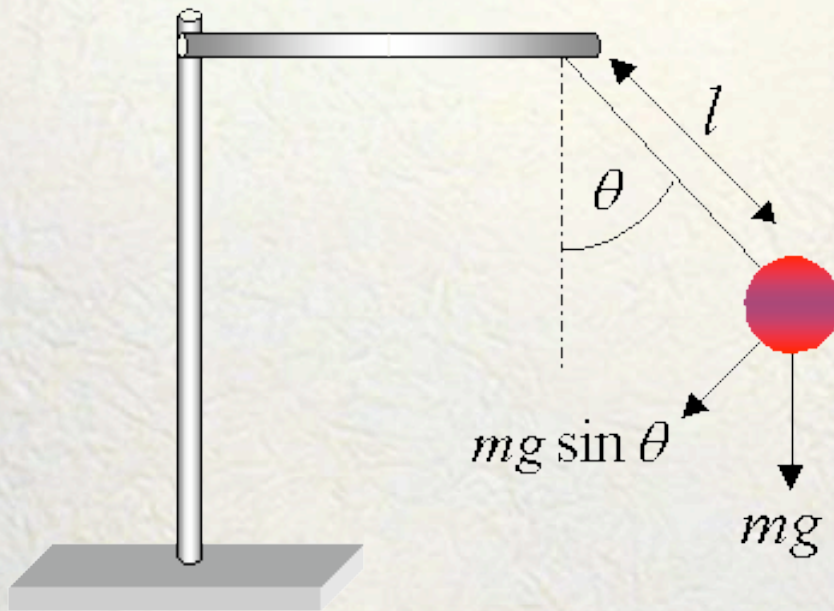
“**Precision** is the ability of an instrument to produce the same value or result, given the same input”

Example: the measurement of the pendulum position



Where does this noise comes from?

The almost “simple” pendulum



Motion equation

$$ml^2 \frac{d^2 \theta}{dt^2} + mgl \sin \theta = 0$$

Small oscillation approximation,
with damping

$$\frac{d^2 \theta}{dt^2} + \left(\frac{1}{q} \right) \frac{d\theta}{dt} + \omega^2 \sin \theta = 0$$

What if we wait long enough ?

The **very small** oscillation limit.
Let the pendulum swing freely...

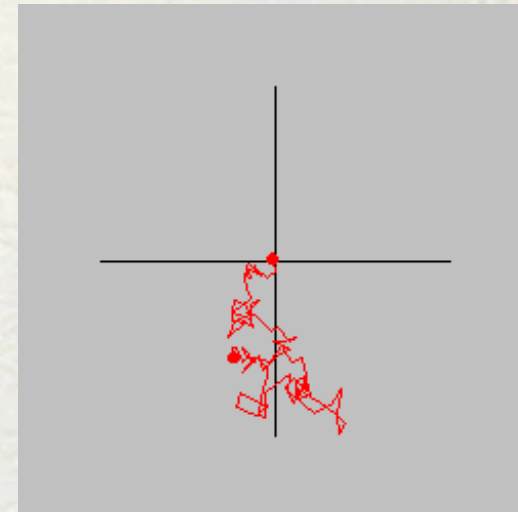


Example: the real measurement of a free swinging pendulum

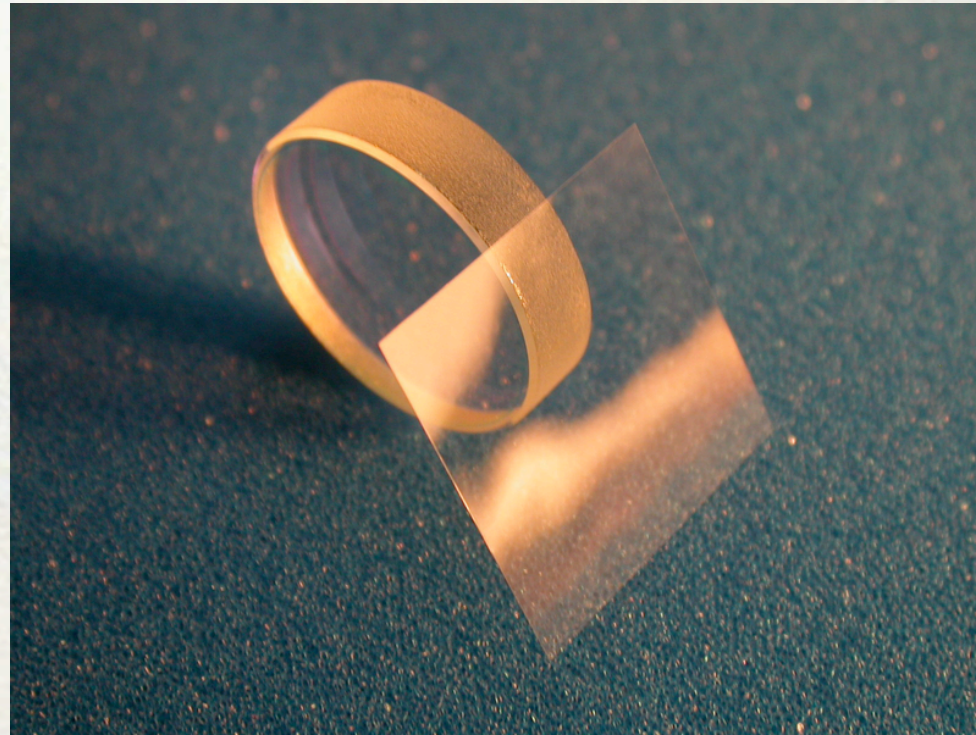
Mass $m = 1 \text{ Kg}$
Length $l = 1 \text{ m}$
rms motion = $2 \cdot 10^{-11} \text{ m}$

Mass $m = 1 \text{ g}$
Length $l = 1 \text{ m}$
rms motion = $6 \cdot 10^{-10} \text{ m}$

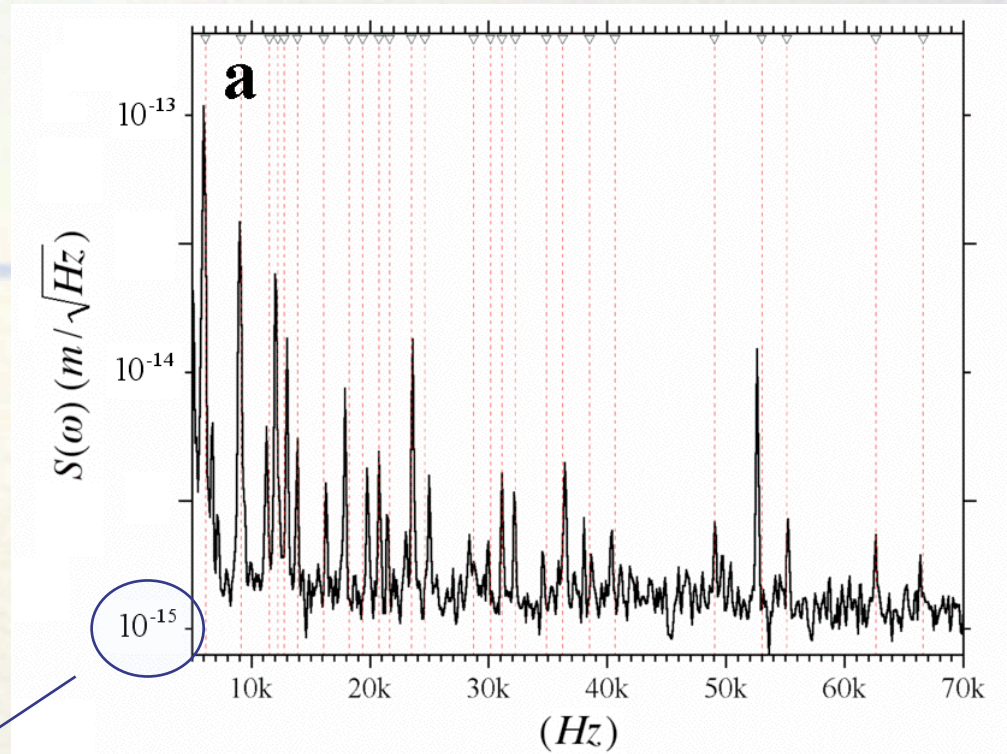
Mass $m = 10^{-6} \text{ g}$
Length $l = 1 \text{ m}$
rms motion approx 1 micron



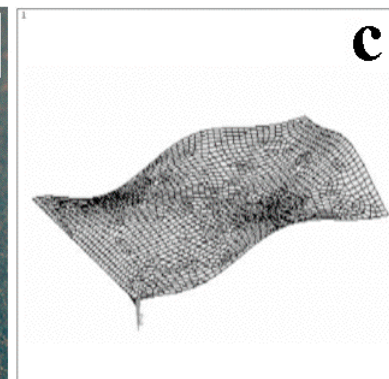
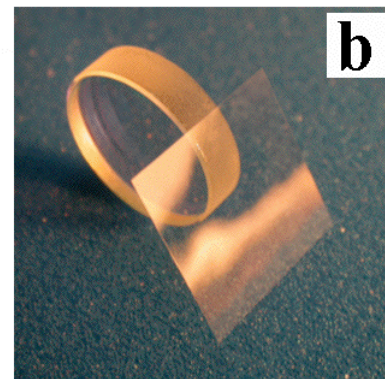
Direct measurement of internal vibration on a thin fused silica slab



Risultati sperimentali



10^{-15}



The mother of all noises: **thermal noise**

Brownian motion → **thermal noise**

thermal noise is the name commonly given to **fluctuations** affecting a physical observable of a macroscopic system at **thermal equilibrium** with its environment.

The **internal energy** of a macroscopic apparatus at thermal equilibrium is shared between all its degrees of freedom or, equivalently, between all its normal modes each carrying an average energy **kT** , where T is the equilibrium temperature.

This is true also for such modes as the oscillations of springs, pendula, needles, etc.

Such an **energy** manifests itself as a **random fluctuation** of the relevant observable experimentally perceived as the noise affecting its measured value.

How to model **thermal noise**



To describe such a dynamics it is necessary to introduce a statistical approach

- a) Fokker-Plank equation
- b) Langevin Equation

The usual equation of motion (Newton $f = ma$) should be changed in order to accomodate non deterministic forces. What we get is called **Langevin equation**:

There are two kinds of forces acting on the pollen grain:

1) A viscous drag (motion in a fluid)

$$-6\pi \eta a \dot{x}$$

2) A fluctuating force (representing the incessant impact of the molecules)

$$\zeta(t)$$

$$m \ddot{x} = -6\pi \eta a \dot{x} + \zeta(t)$$

Generalization:

The viscous drag expression can be generalized in order to describe a wider class of damping functions

$$-m \int_{-\infty}^t \gamma(t - \tau) \dot{x} d\tau$$

Generalized Langevin equation

$$m \ddot{x} = -m \int_{-\infty}^t \gamma(t - \tau) \dot{x} d\tau + \zeta(t)$$

$\zeta(t)$ Is the stochastic force with known statistical properties:

Probability density function, moments, correlations, ...

For the pendulum we get:

$$\ddot{x} = -\omega_p^2 x - \int_{-\infty}^t \gamma(t-\tau) \dot{x} d\tau + \frac{\xi(t)}{m}$$

The arbitrariness in our choice of the damping function and the stochastic force is constrained by the so-called **Fluctuation-Dissipation theorem** which suffices to enforce the energy equipartition

$$\langle \xi(t) \xi(0) \rangle = k T m \gamma(|t|)$$

In the spectral domain, for a linear system, is always possible to write its response to an external force like:

$$X(\omega) = H(\omega)F(\omega)$$

Where H is the system transfer function.

$$H(\omega) = H'(\omega) + i H''(\omega) = |H(\omega)| e^{i\phi(\omega)}$$

The F-D Theorem can be written here as:

$$S_x(\omega) = -4 k T \frac{H''(\omega)}{\omega}$$

The **dissipative properties** of the dynamical system are thus directly related to the equilibrium fluctuations.

Fluctuation-Dissipation theorem

The **dissipative properties** of the dynamical system are thus directly related to the equilibrium fluctuations.

Physical connection:

the source of the **fluctuations** is the very same of the source of the **dissipation**

Why should we bother with **thermal noise** ?



**Physics is like sex.
Sure, it may give some
practical results,
but that's not why we do
it.**

Richard Feynman

Cambio di prospettiva

Anziché considerare il rumore un mero disturbo, ipotizzare un suo impiego costruttivo.

- Due esempi:
- 1) La **Risonanza Stocastica**
 - 2) Effetto **Wisepower**

Dynamical description

For the generic physical dynamic system we can use the “potential” description (conservative forces):

$$\ddot{x} = -V'(x) - f(\dot{x}) + \zeta(t)$$

Where we set $m=1$ and introduced the dissipation function f

For the linear pendulum = harmonic oscillator

$$V(x) = \frac{1}{2} m \omega_p^2 x^2$$

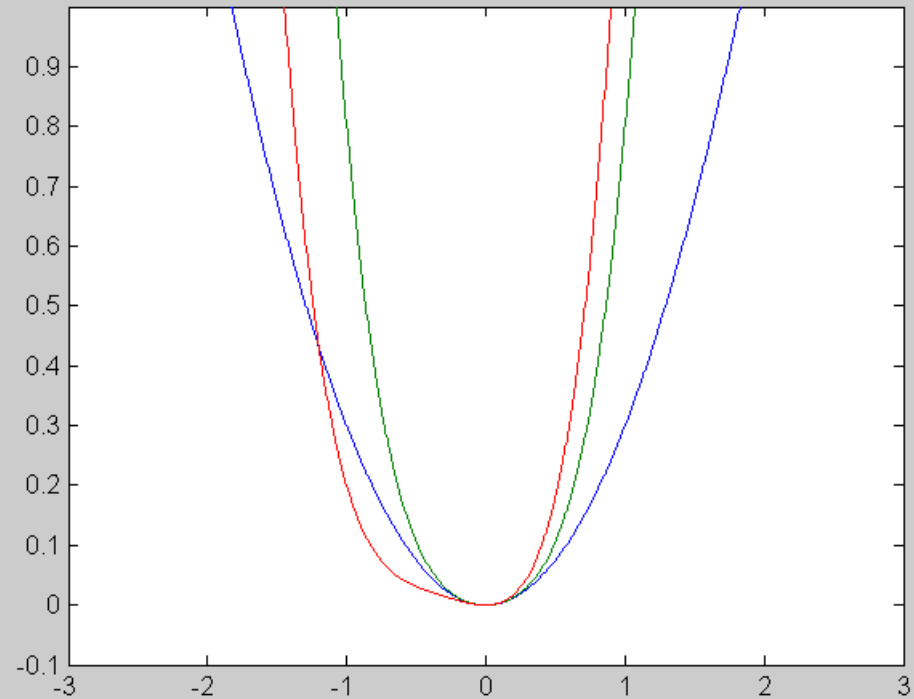
The potential of a linear system is a parabolic one!

Langevin Equation for the linear pendulum

$$\ddot{x} = -\omega_p^2 x - \int_{-\infty}^t \gamma(t - \tau) \dot{x} d\tau + \frac{\zeta(t)}{m}$$

nonlinear dynamics in **nonlinear** potential

We can
change the
shape of the
curve by
adding terms
to the
polynomial
expression of
the potential



Monostable nonlinear potential

Bistable (nonlinear) potential

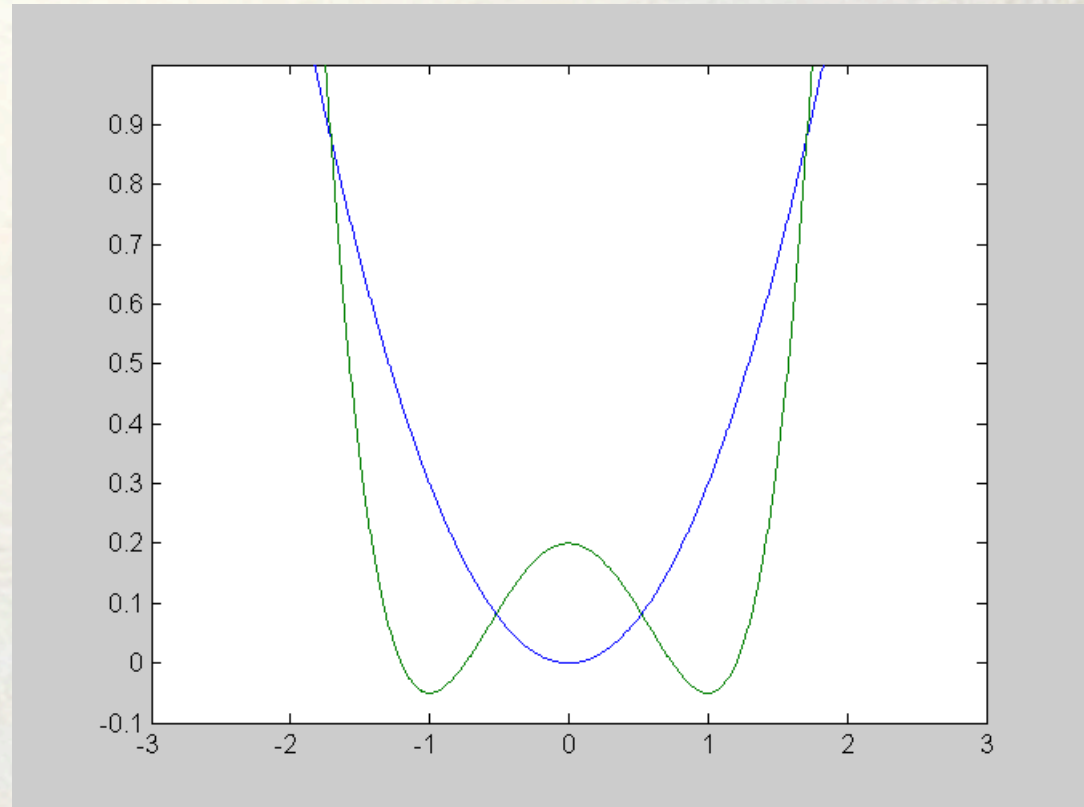
$$V(x) = -\frac{1}{2}ax^2 + \frac{1}{4}bx^4$$

Bistable ...

Add cubic term

Add quartic term

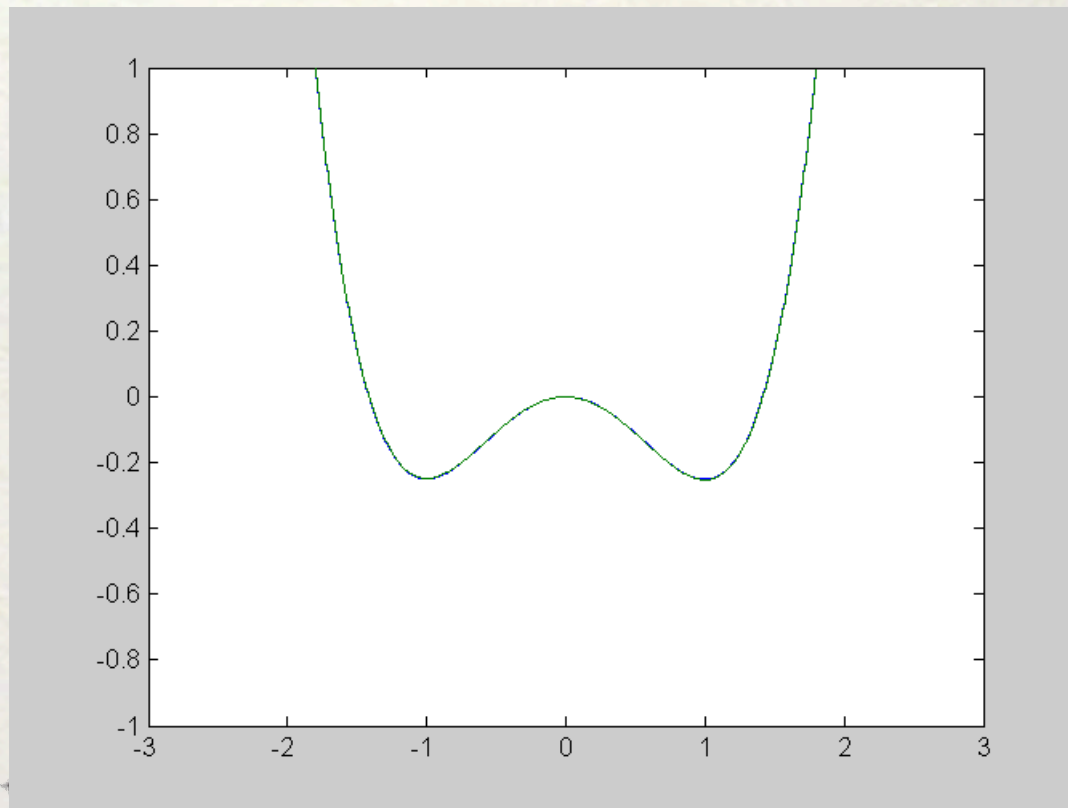
symmetric
quartic



What if ...

I introduce the time dependent part into the potential ?

$$V(x,t) = V(x) + A(t) \quad A(t) = A \sin(\omega t)$$

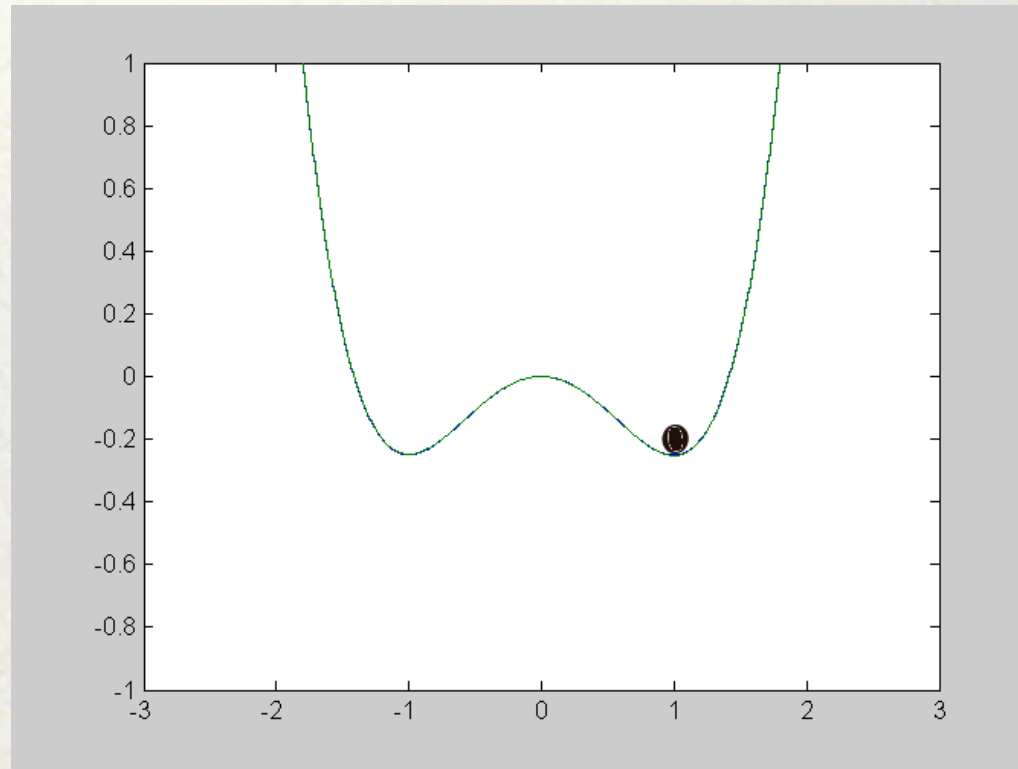




Marble in a cartoon: let's play dynamics!

$$\ddot{x} = -V'(x) - \gamma\dot{x} + A(t)$$

$$V(x) = -\frac{1}{2}ax^2 + \frac{1}{4}bx^4 \quad A(t) = A \sin(\omega t)$$



The **Stochastic Resonance** phenomenon

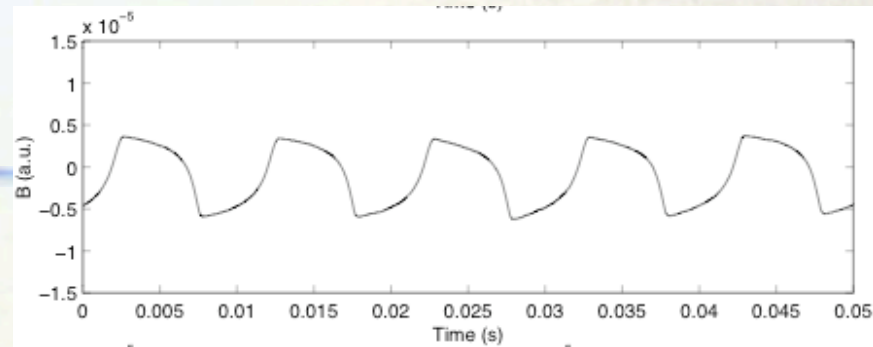
Let's consider our very basic ingredients:

$$\ddot{x} = -V'(x) - \gamma\dot{x} + A(t)$$

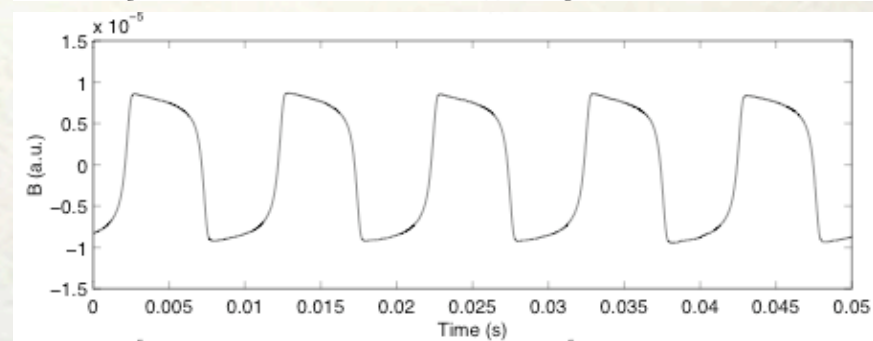
$$V(x) = -\frac{1}{2}ax^2 + \frac{1}{4}bx^4 \quad A(t) = A \sin(\omega t)$$

Let's look for the time evolution of $x(t)$

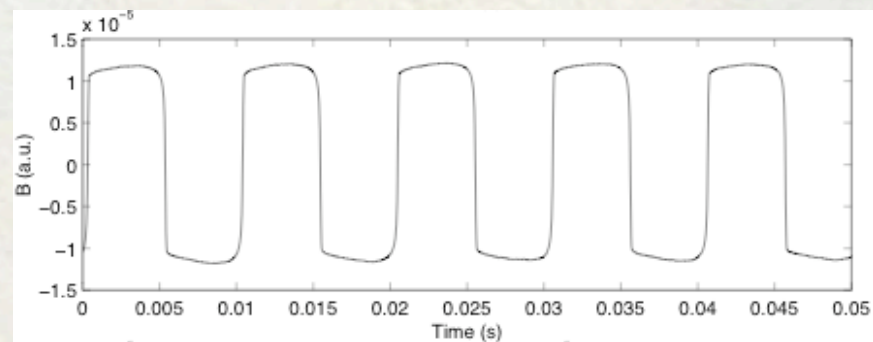
How does it look?



$A = 3$

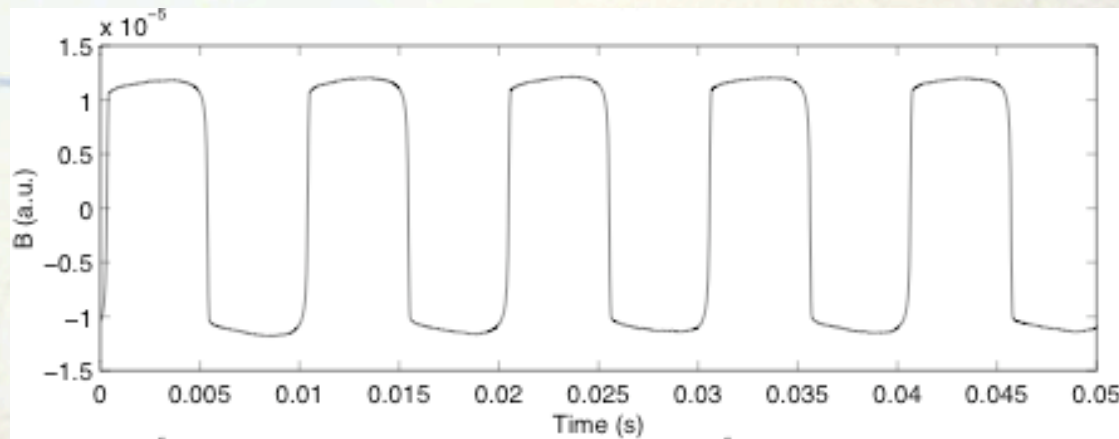


$A = 4$

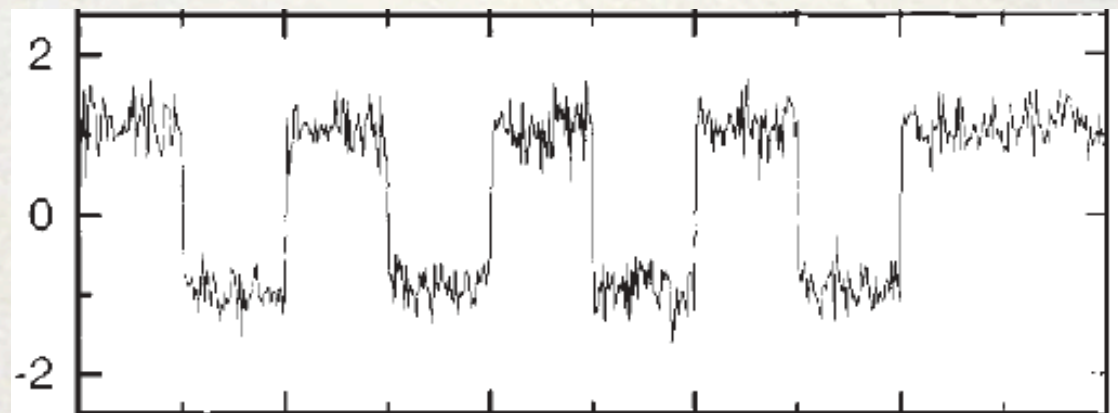


$A = 5$

What if we add the noisy force?



No noise

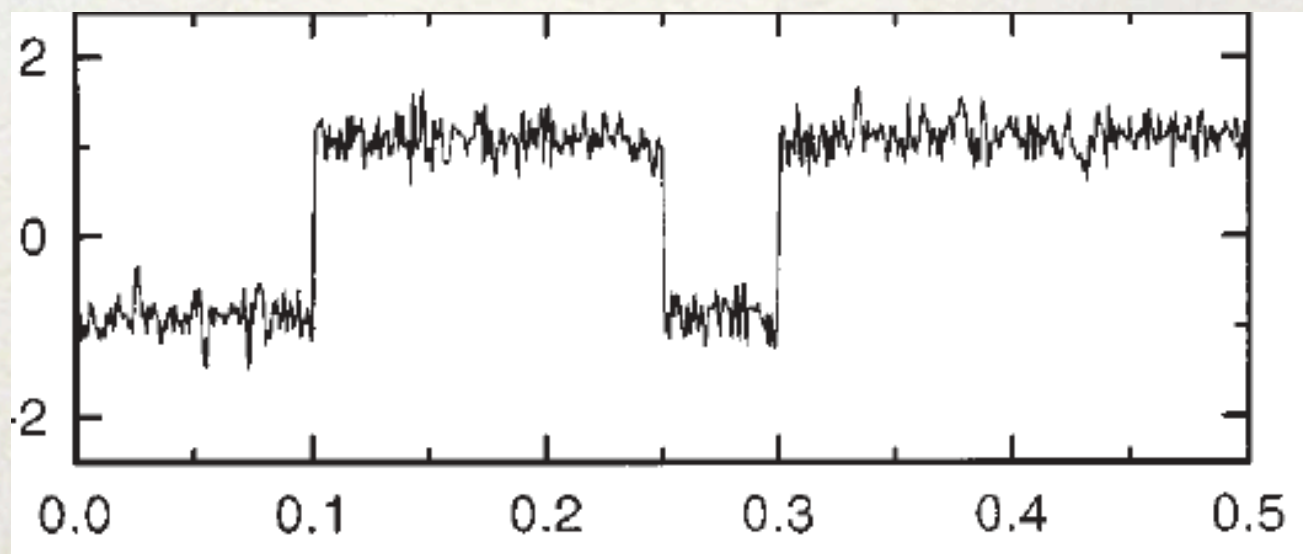


With noise

What if the periodic driving is too small
to let the marble to jump??

The marble would stay trapped into one well unless
The noise provide some help...

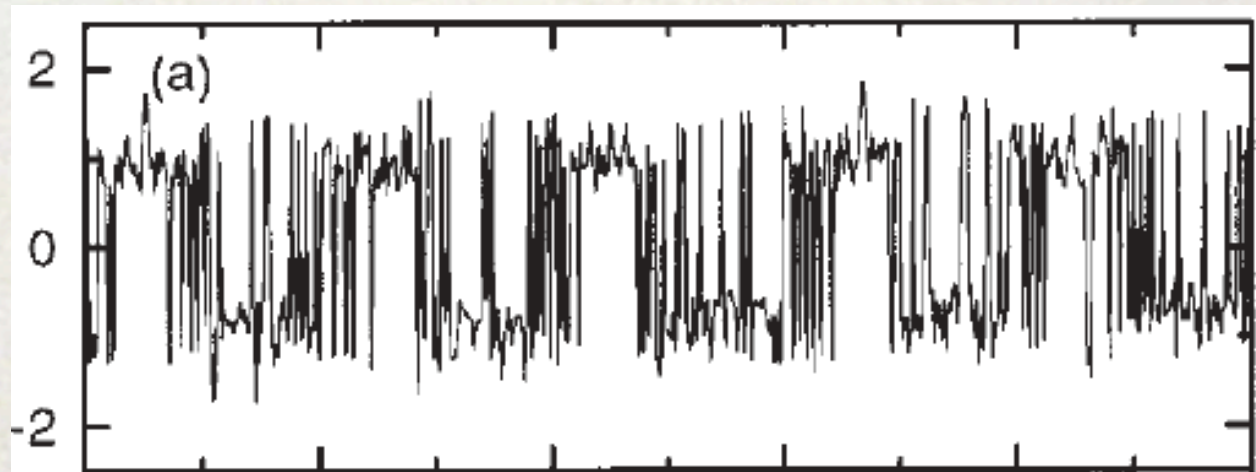
For **small** noise we will have **occasional** jumps



What if the periodic driving is too small
to let the marble to jump??

The marble would stay trapped into one well unless
The noise provide some help...

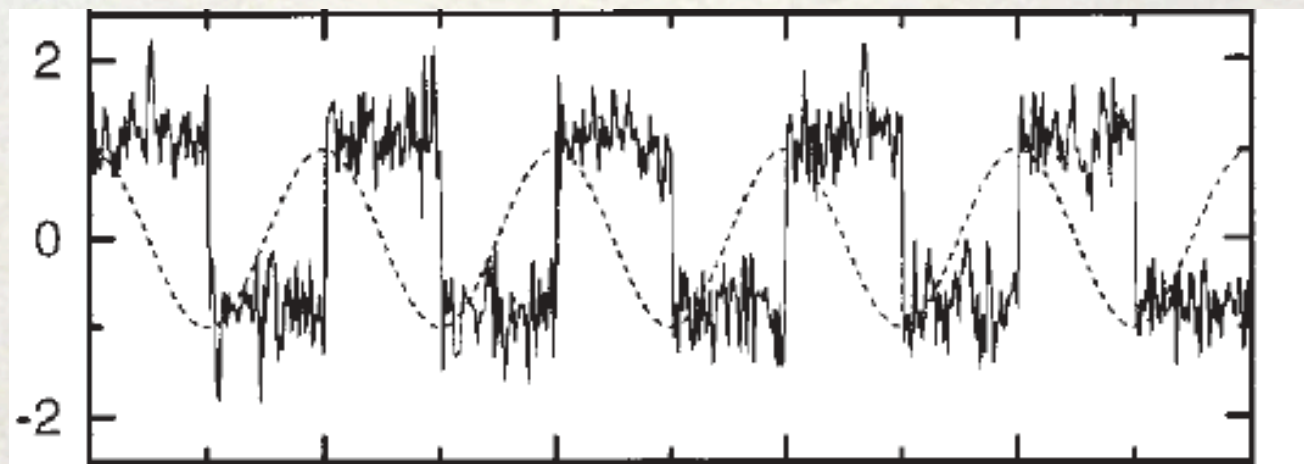
For **large** noise we will have **frequent** jumps



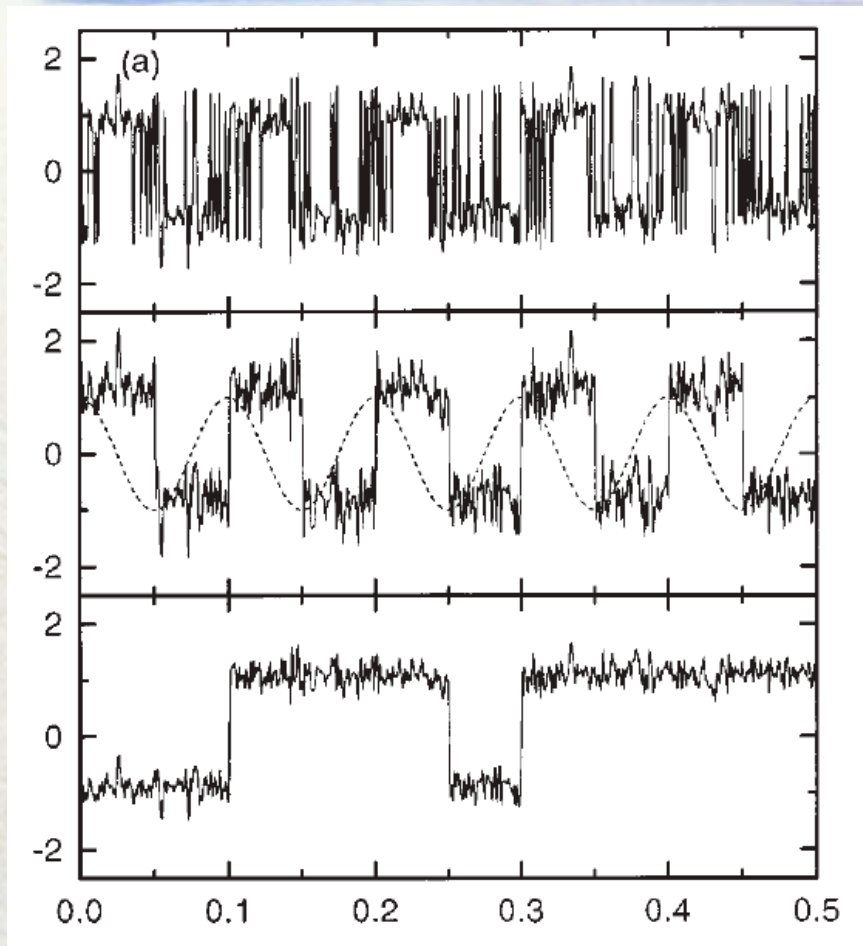
What if the periodic driving is too small
to let the marble to jump??

And in between ?

It exists an optimal noise intensity that produces the
Jumps **in synchrony** with the forcing !!!



The **Stochastic Resonance** Phenomenon



Large noise

Optimal noise

Small noise

The **Stochastic Resonance** Phenomenon

First paper in 1981 by Benzi, Parisi, Sutera, Vulpiani.

Since then more than 3000 papers (to date)...

stochastic Resonance
1998**SR**2008

The European Physical Journal B

Vol. 69 No. 1 (May I 2009)

Special Issue: Stochastic Resonance

Perugia, Aug. 17-21, 2008

THE CONFERENCE

Other “surprising” noise induced phenomena

1. Resonant Trapping
2. Dithering Effect
3. Resonant Activation
4. Brownian Ratchet
5. Resonant crossing
6. ...
7. **Energy harvesting (noise harvesting)**

Energy harvesting basic ideas

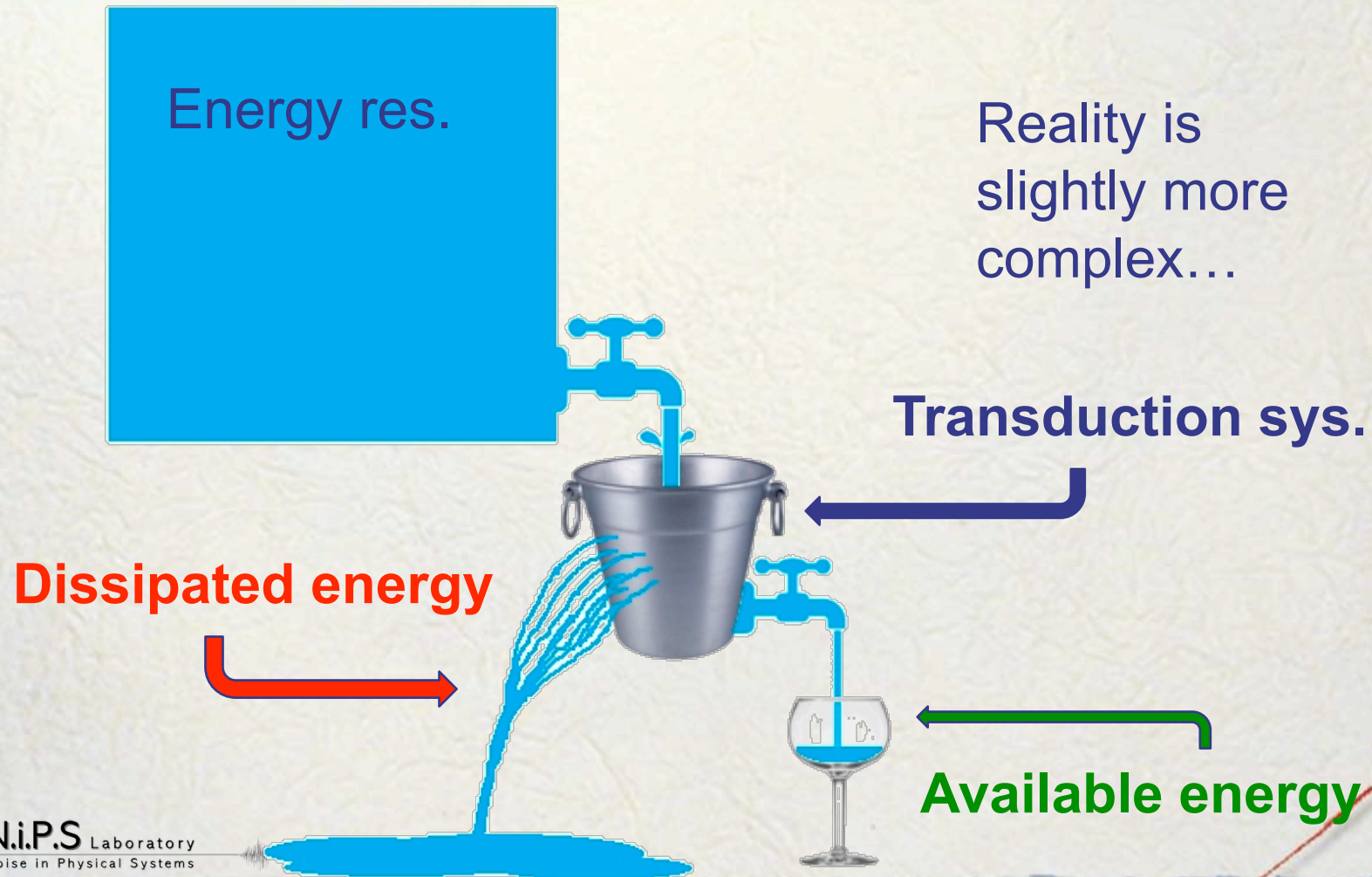


ENERGY



Unlimited source of free energy, readily available for multiple uses...

Energy harvesting basic ideas



Energy harvesting basic ideas

Kinetic energy

wind

sound

Falling bodies

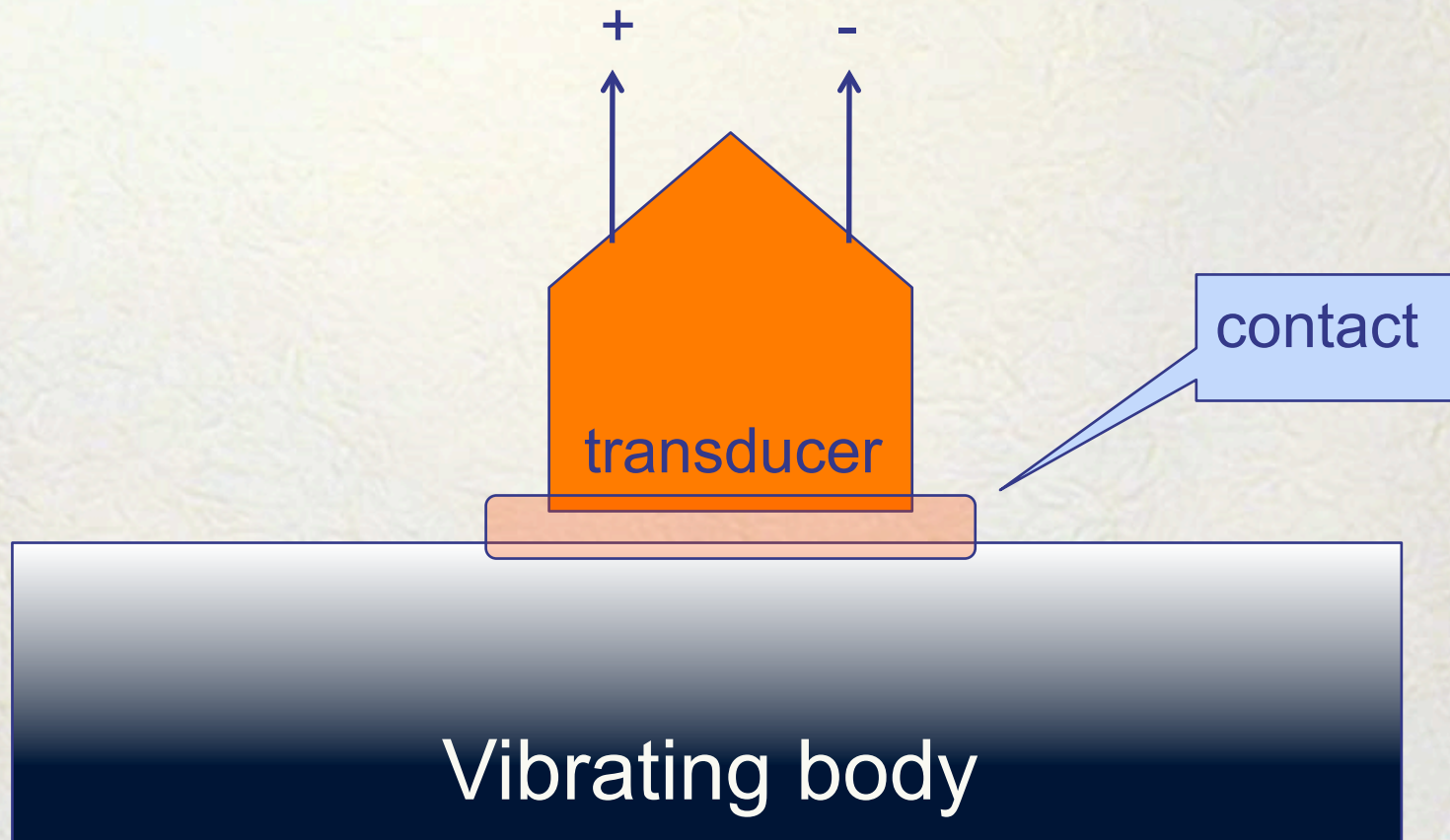
vibrations

water waves and tides

Focus on vibrations of solid bodies....

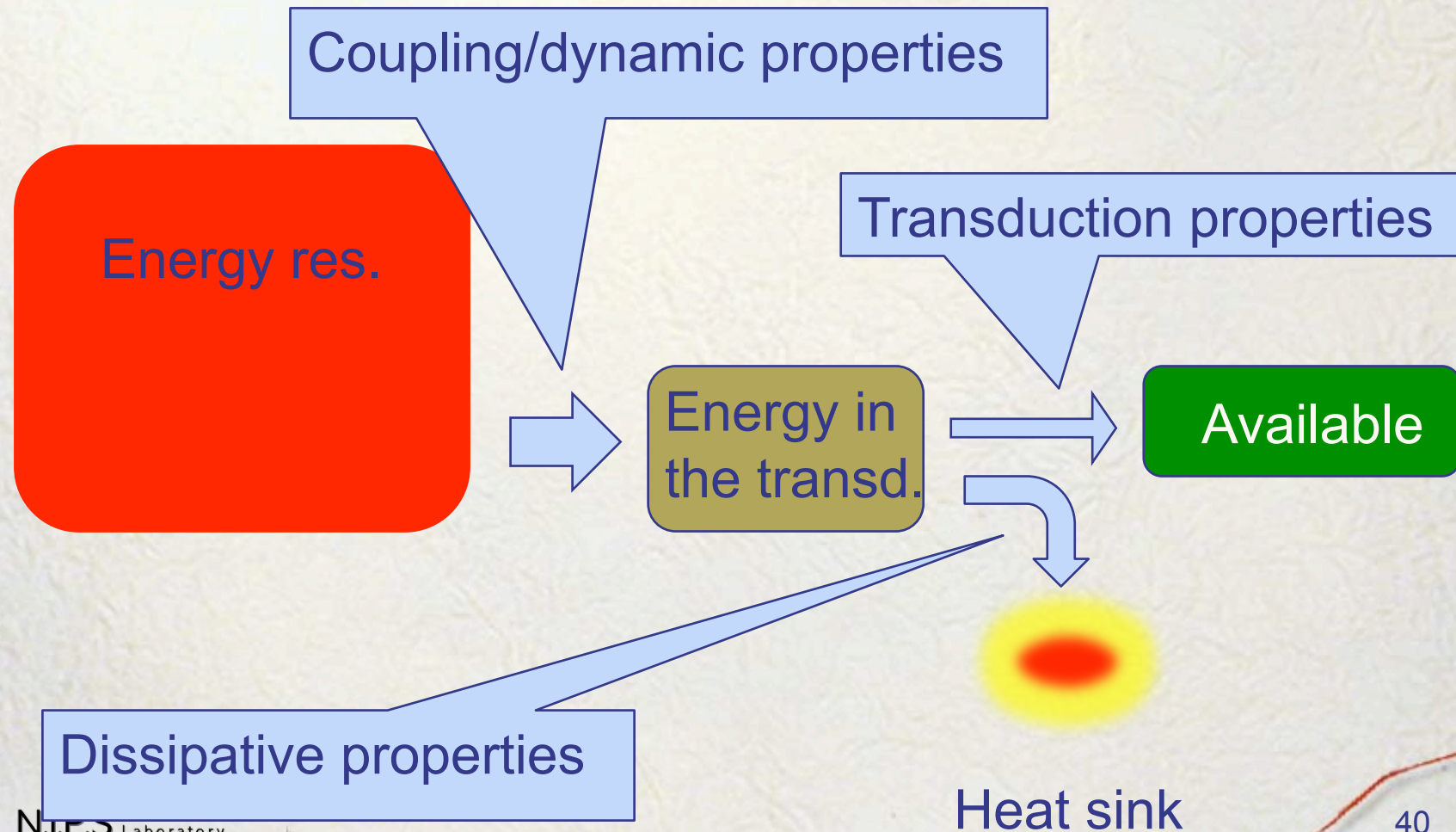
Vibrations energy harvesting

Basic Scheme



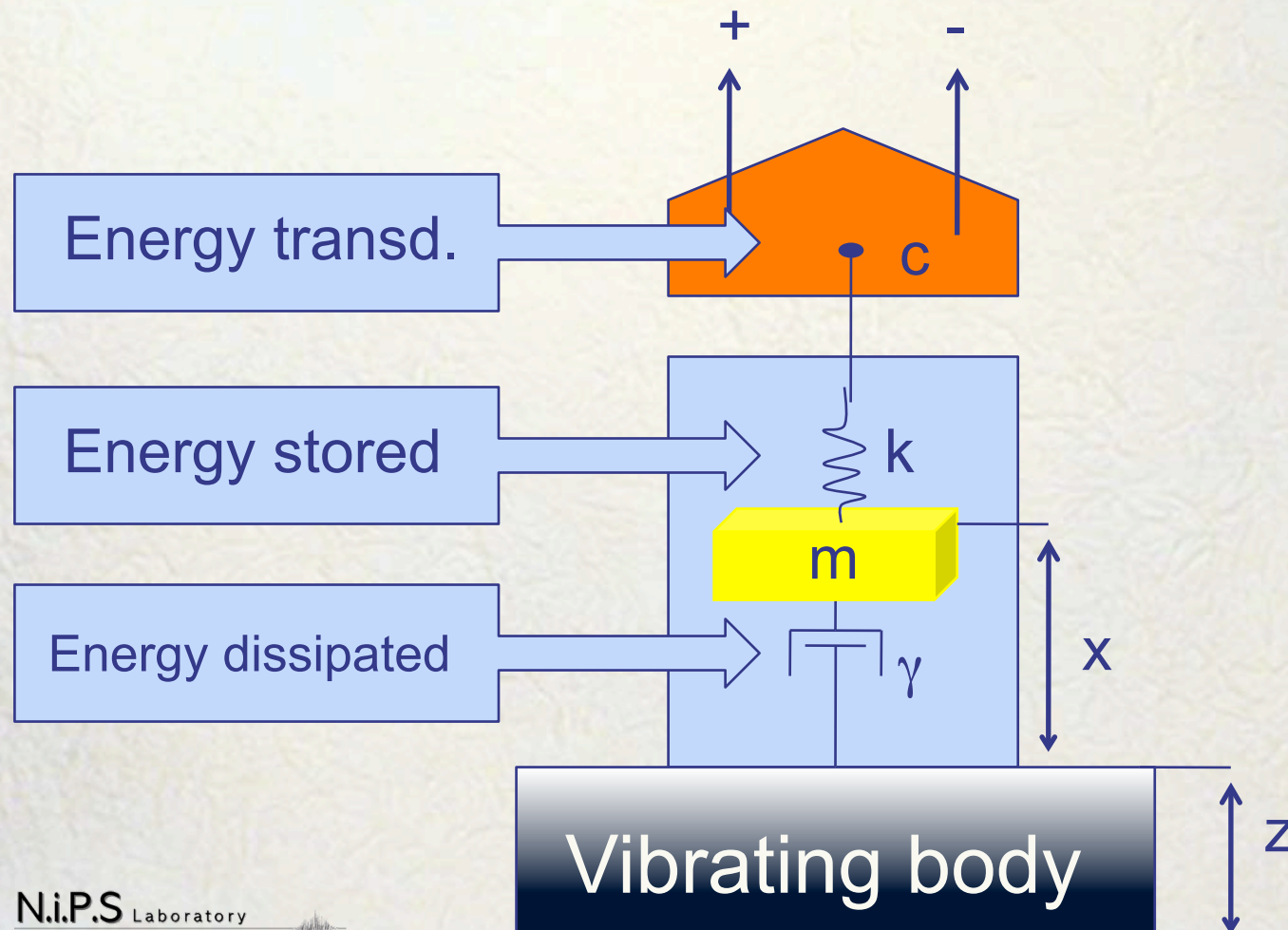
Vibrations energy harvesting

Energy budget

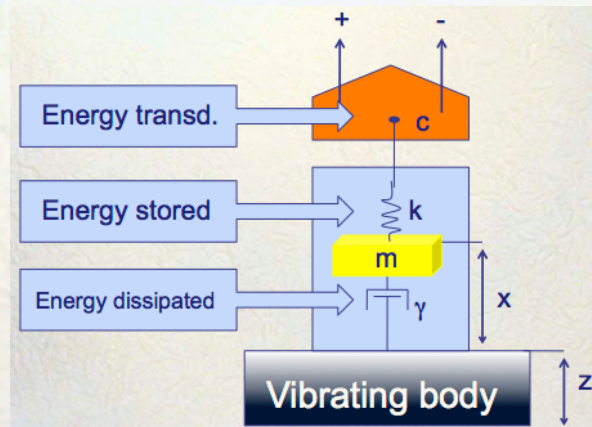


Vibrations energy harvesting

Dynamical model



Vibrations energy harvesting

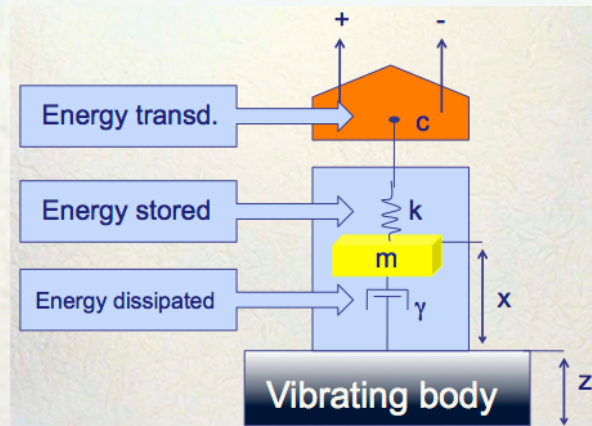


Dynamical model

$$m\ddot{x} = -\frac{dU(x)}{dx} - \gamma\dot{x} - c(x,V) + \xi_z$$

Where: $U(x)$ Represents the Energy stored
 $\gamma\dot{x}$ Accounts for the Energy dissipated
 $c(x,V)$ Accounts for the Energy transduced
 ξ_z Accounts for the input Energy

Vibrations energy harvesting



Dynamical model

$$\left\{ \begin{array}{l} m\ddot{x} = -\frac{dU(x)}{dx} - \gamma\dot{x} - c(x,V) + \xi_z \\ \dot{V} = F(\dot{x},V) \end{array} \right.$$

Equations that link the vibration-induced displacement
with the Voltage

Vibrations energy harvesting

Dynamical model

$$\left\{ \begin{array}{l} m\ddot{x} = -\frac{dU(x)}{dx} - \gamma\dot{x} - c(x,V) + \xi_z \\ \dot{V} = F(\dot{x},V) \end{array} \right.$$

Details depend on the physics...

The diagram shows a set of two equations enclosed in a large left-facing curly bracket. The first equation is $m\ddot{x} = -\frac{dU(x)}{dx} - \gamma\dot{x} - c(x,V) + \xi_z$. The second equation is $\dot{V} = F(\dot{x},V)$. The term $F(\dot{x},V)$ is enclosed in a blue circle. A red rounded rectangle containing the text "Details depend on the physics..." has an arrow pointing to the blue circle. Another arrow points from the blue circle to the term $c(x,V)$ in the first equation.

Three main transduction mechanisms...

Vibrations energy harvesting

Transduction mechanisms

1

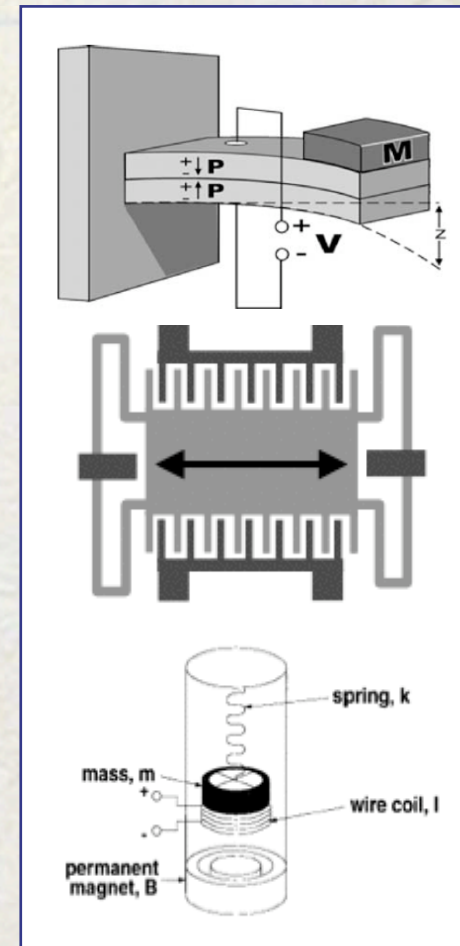
Piezoelectric: dynamical strain is converted into voltage difference.

2

Capacitive: geometrical variations induce voltage difference

3

Inductive: dynamical oscillations of magnets induce electric current in coils

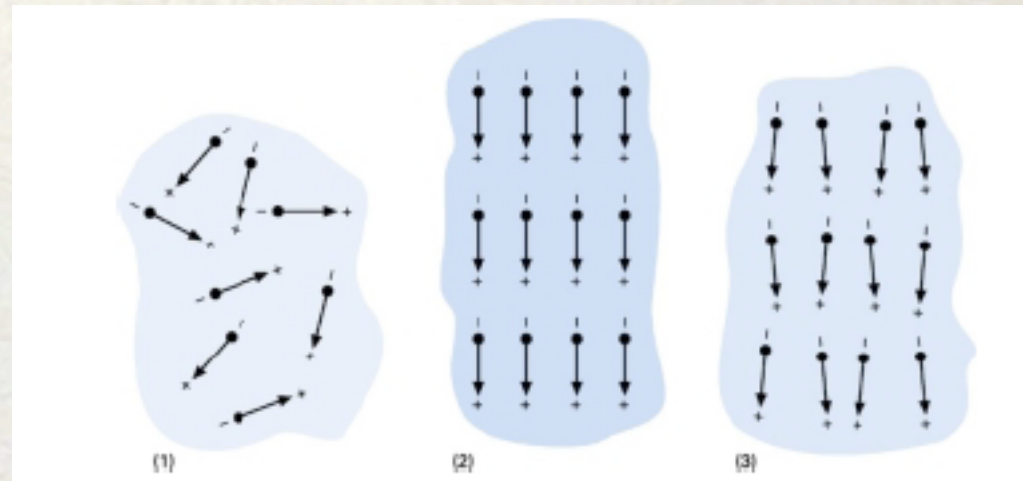
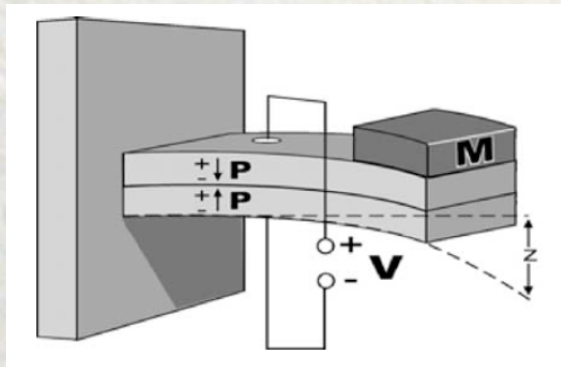


Vibrations energy harvesting

Transduction mechanisms

1

Piezoelectric: dynamical strain is converted into voltage difference.



Vibrations energy harvesting

Transduction mechanisms

1

Piezoelectric: dynamical strain is converted into voltage difference.

$$m\ddot{x} = -\frac{dU(x)}{dx} - \gamma\dot{x} - K_V V + \xi_z$$

$$\dot{V} = K_c \dot{x} - \frac{1}{\tau_p} V$$

The Physics of piezo materials

The available power is proportional to V^2

Vibrations energy harvesting

Transduction mechanisms

$$\left\{ \begin{array}{l} m\ddot{x} = -\frac{dU(x)}{dx} - \gamma\dot{x} - K_v V + \xi_z \\ \dot{V} = K_c \dot{x} - \frac{1}{\tau_p} V \end{array} \right.$$

$U(x)$ Represents the Energy stored

When $U(x) = \frac{1}{2} kx^2$ it is called a linear system

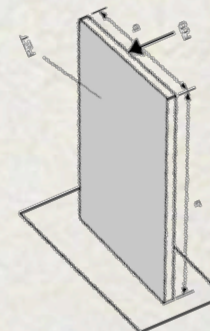
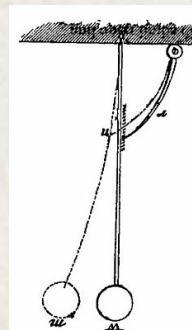
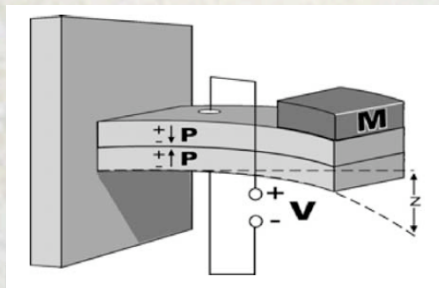
Vibrations energy harvesting

Linear systems

When $U(x) = \frac{1}{2} kx^2$ it is called a linear system

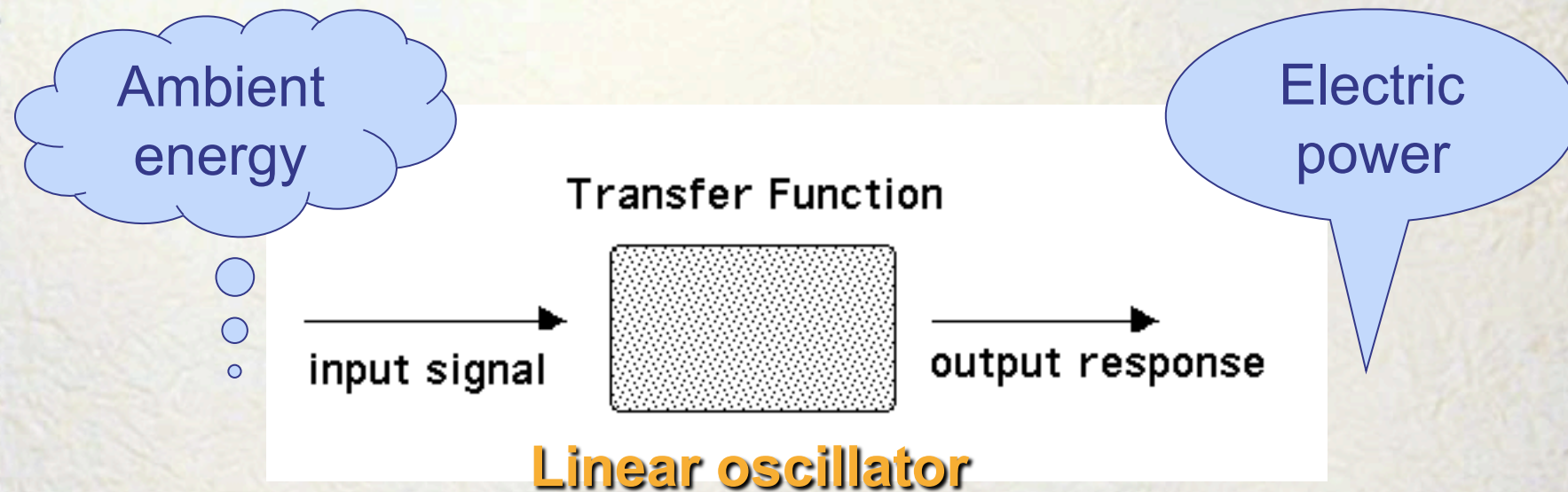
Linear systems have some interesting features... (and engineers like them most)

- 1) There exist a simple math theory to solve the eq.s
- 2) They have a resonant behaviour (resonance freq.)
- 3) They can be “easily” realized with catilevers and pendula



Vibrations energy harvesting

Linear systems

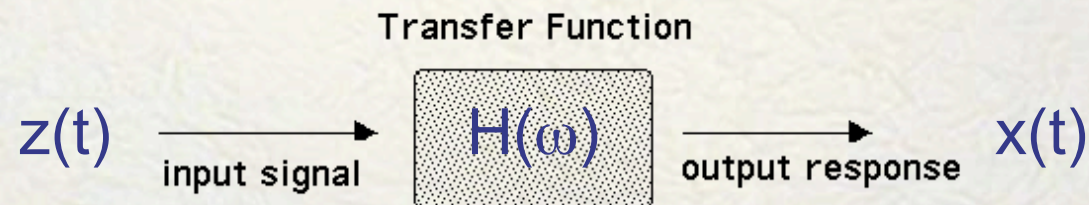


The transfer function is a math function of the frequency, in the complex domain, that can be used to represent the performance of a linear system

Vibrations energy harvesting

Linear systems

In a linear system, thanks to the transfer function $H(\omega)$, the output spectrum can be obtained from the input spectrum through a simple multiplication...



$$S_x(\omega) = |H(\omega)|^2 S_z(\omega)$$

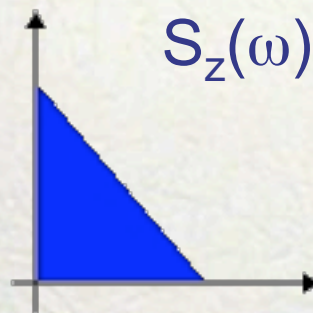
Output power
spectrum

Input power
spectrum

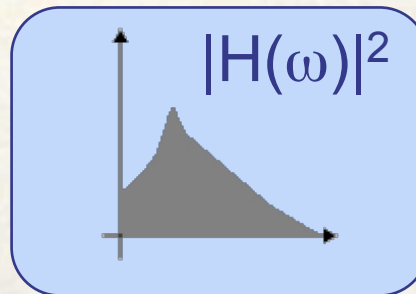
Vibrations energy harvesting

Linear systems

The transfer function is important because it acts as a filter on the incoming energy...

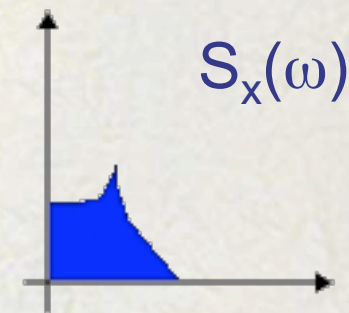


Freq. spectrum
of the available
energy



Transfer
function of the
transducer

$$S_x(\omega) = |H(\omega)|^2 S_z(\omega)$$



Freq. spectrum
of the usable
energy

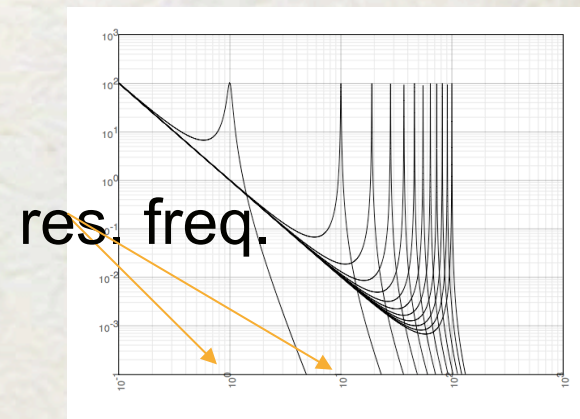
Vibrations energy harvesting

Linear systems

For a linear system the transfer function presents one or more peaks corresponding to the resonance frequencies and thus it is efficient mainly when the incoming energy is abundant in that regions...

This is a serious limitation when you want to build a small energy harvesting system...

Why ?!!



Vibrations energy harvesting

Linear systems

For two main reasons...

- (1) the frequency spectrum of available vibrations instead of being sharply peaked at some frequency is usually very broad.
- (2) The frequency spectrum of available vibrations is particularly rich in energy in the low frequency part... and it is very difficult, if not impossible, to build small low-frequency resonant systems...

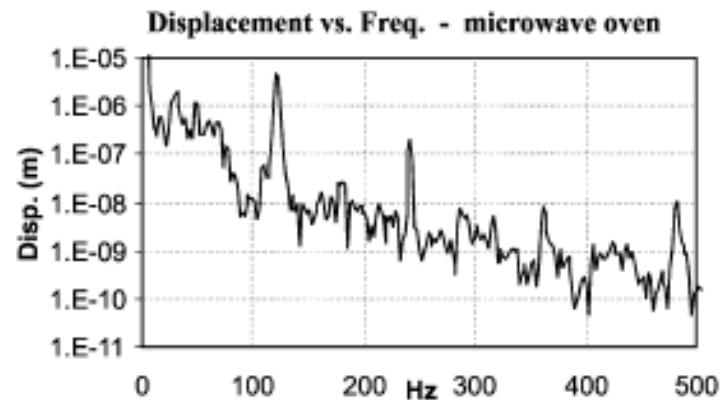
Let's see some examples...



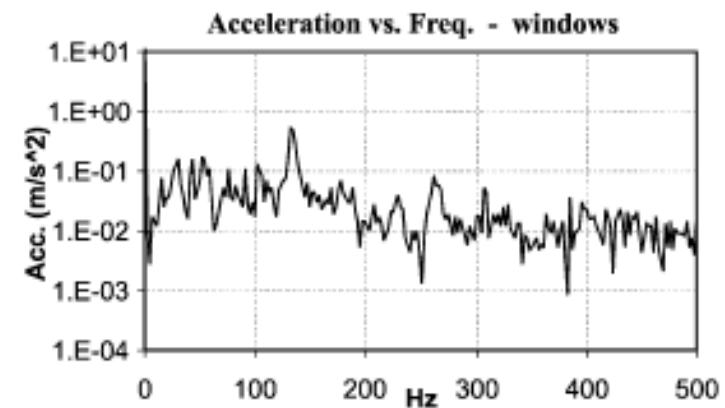
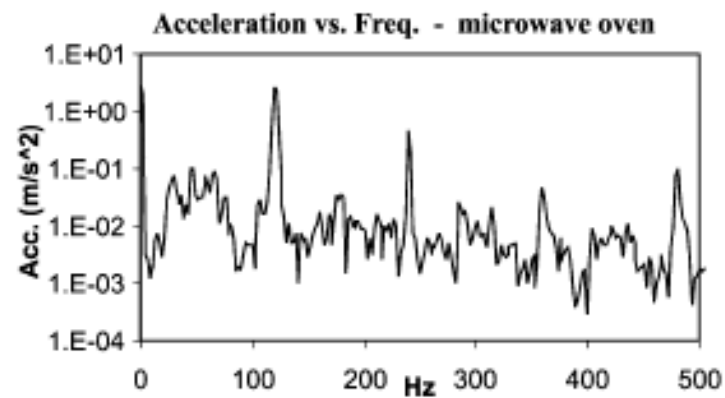
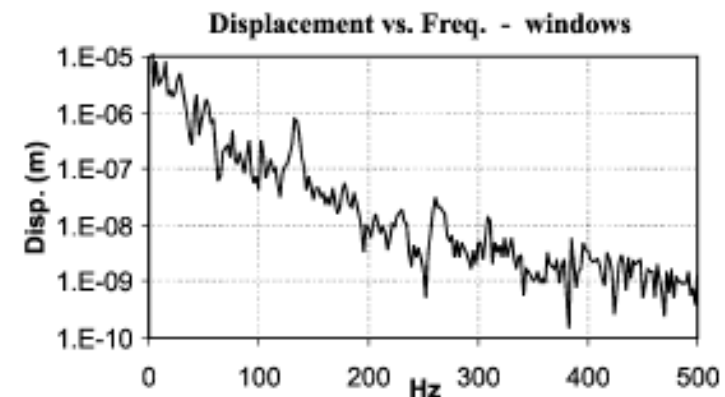
Vibrations energy harvesting

S. Roundy et al. / Computer Communications 26 (2003) 1131–1144

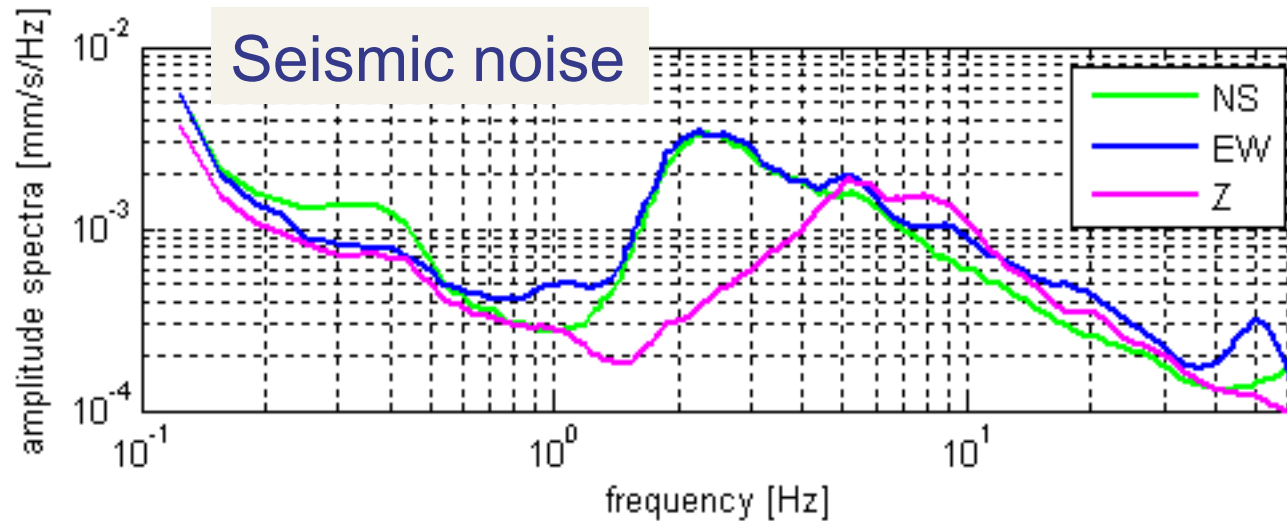
Microwave Casing



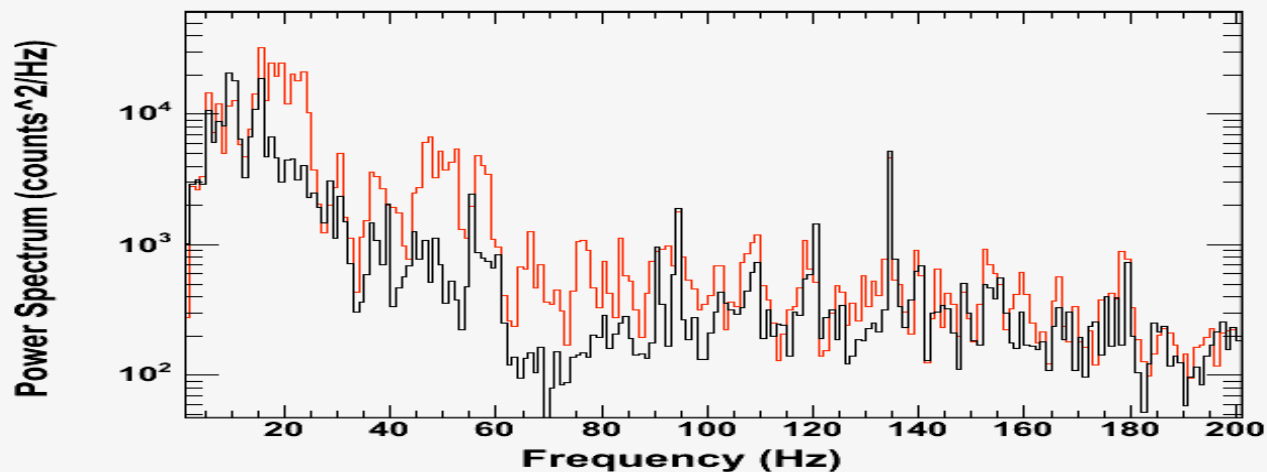
Windows Next to a Busy Street



Vibrations energy harvesting

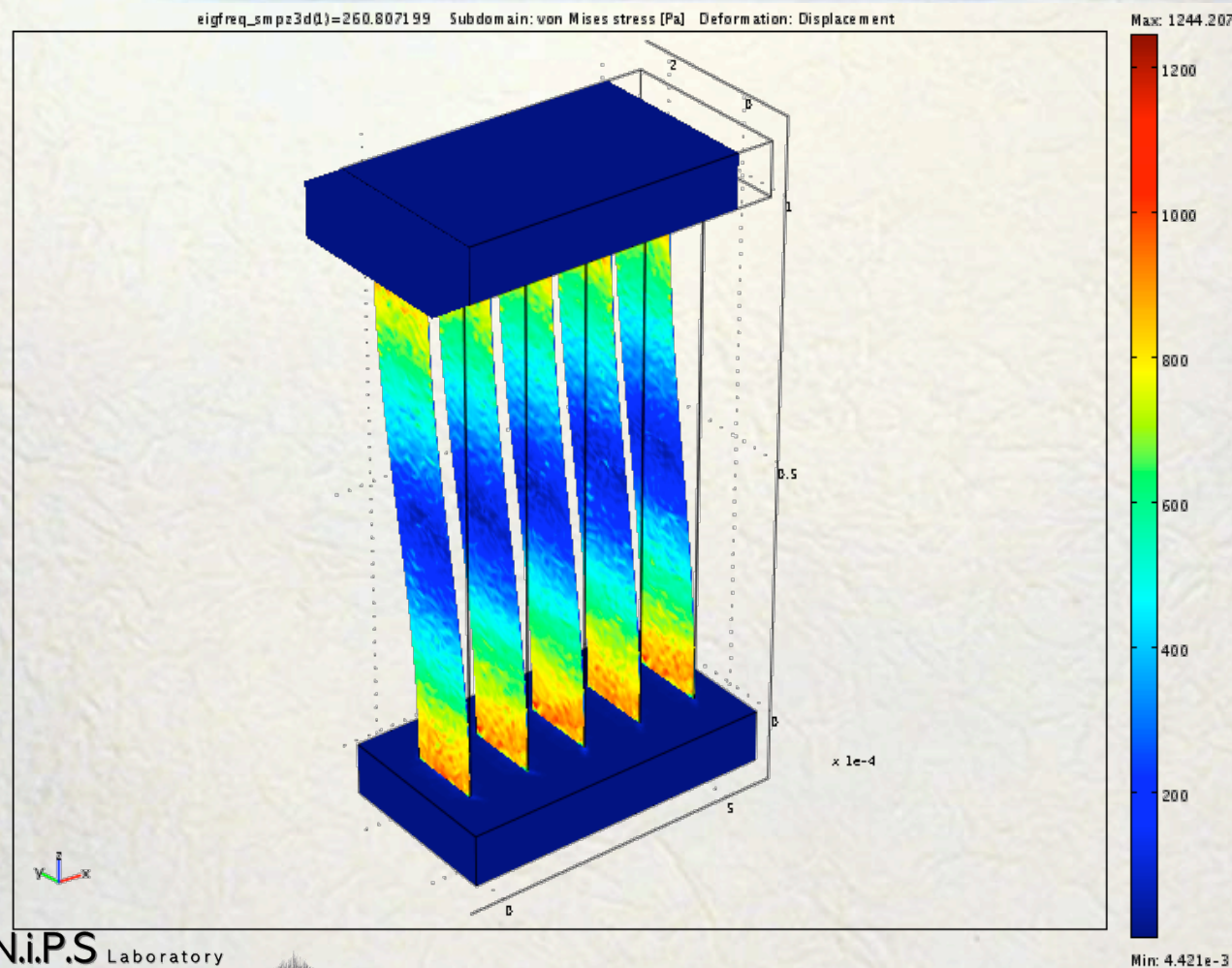


Acoustic noise – quiete working env.



Vibrations energy harvesting

Micro energy harvesting system...



25 μm thick
1 mm high

Freq. 10 KHz

Vibrations energy harvesting

Whish list for the perfect vibration harvester

- 1) Capable of harvesting energy on a broad-band
- 2) No need for frequency tuning
- 3) Capable of harvesting energy at low frequency



- 1) Non-resonant system
- 2) “Transfer function” with wide frequency resp.
- 3) Low frequency operated



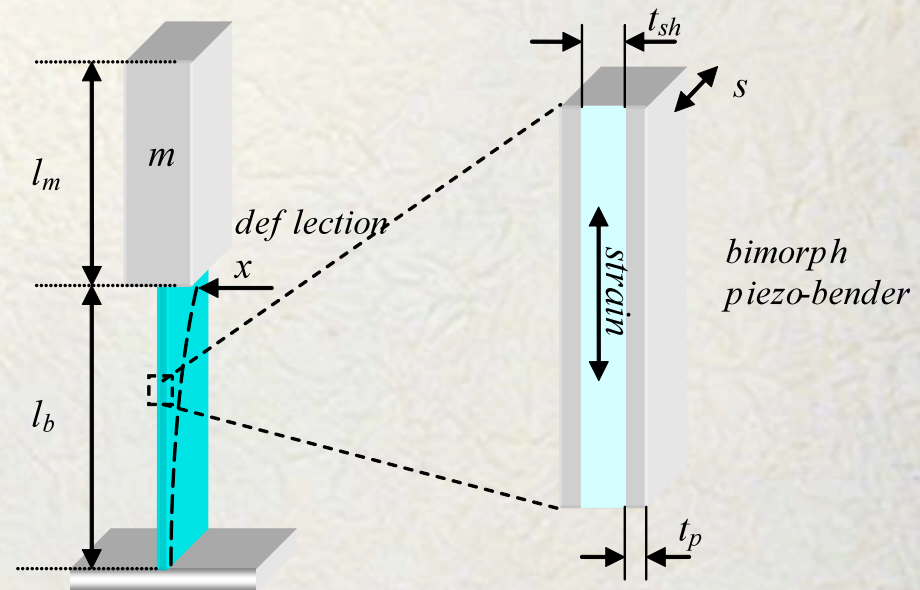
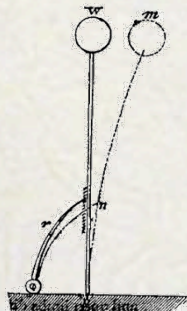
Vibrations energy harvesting

NON-Linear mechanical oscillators

- 1) Non-resonant system
- 2) “Transfer function” with wide frequency resp.
- 3) Low frequency operated

Example...

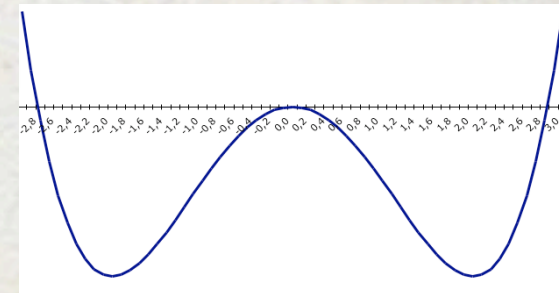
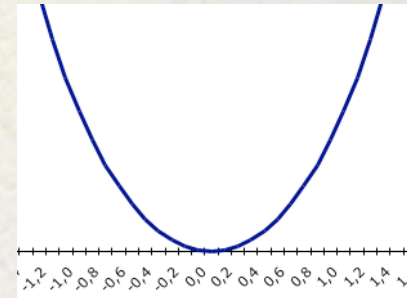
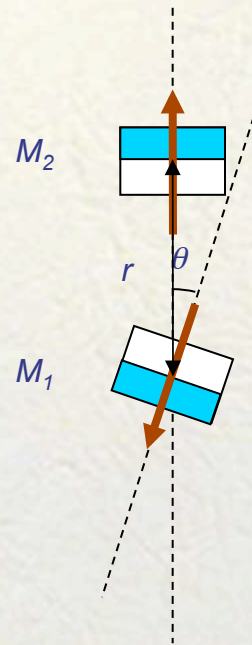
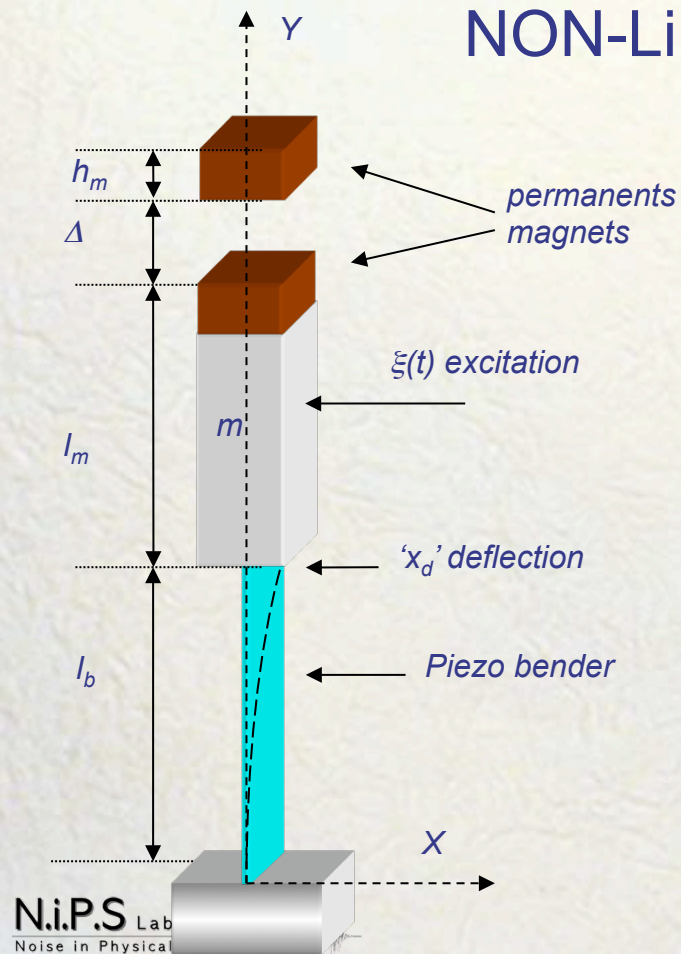
Inverted pendulum



Vibrations energy harvesting

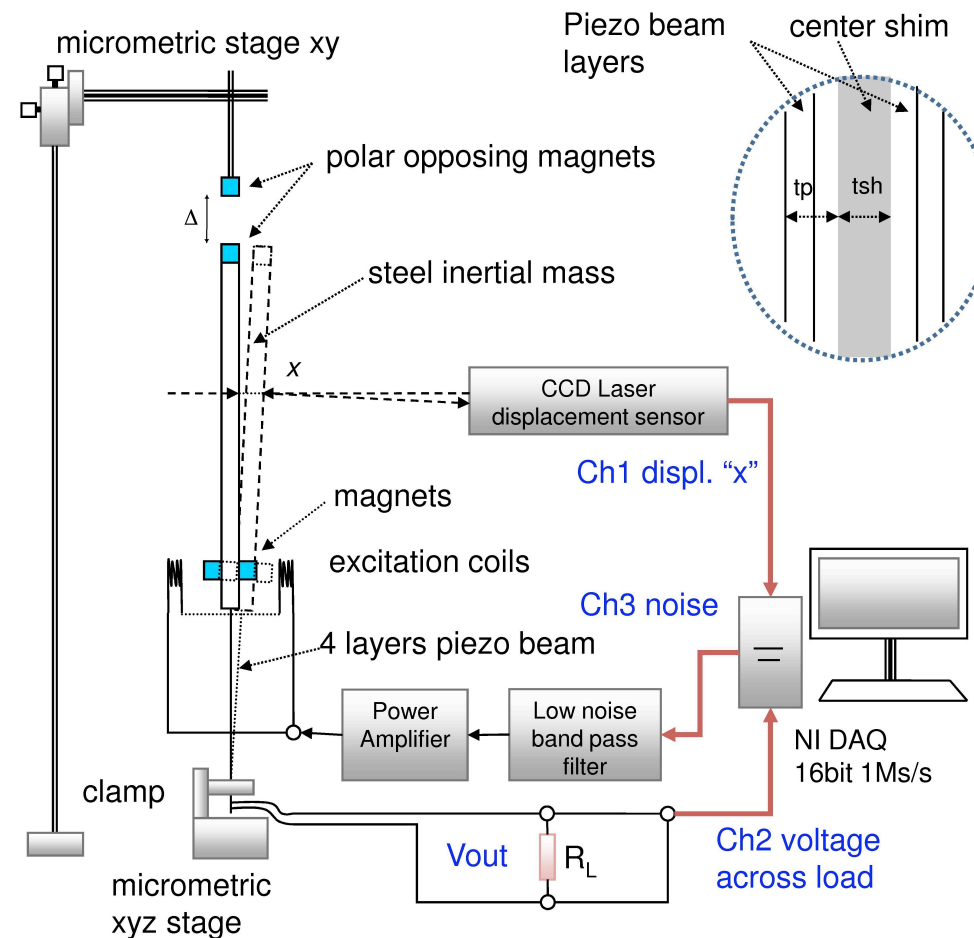
NON-Linear mechanical oscillators

NON-Linear Inverted pendulum



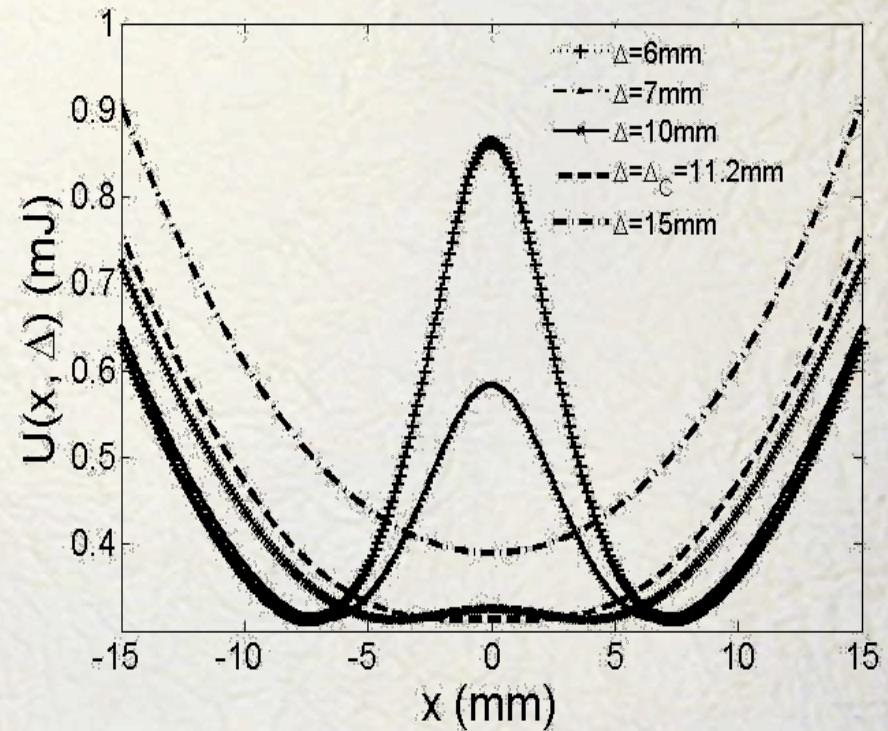
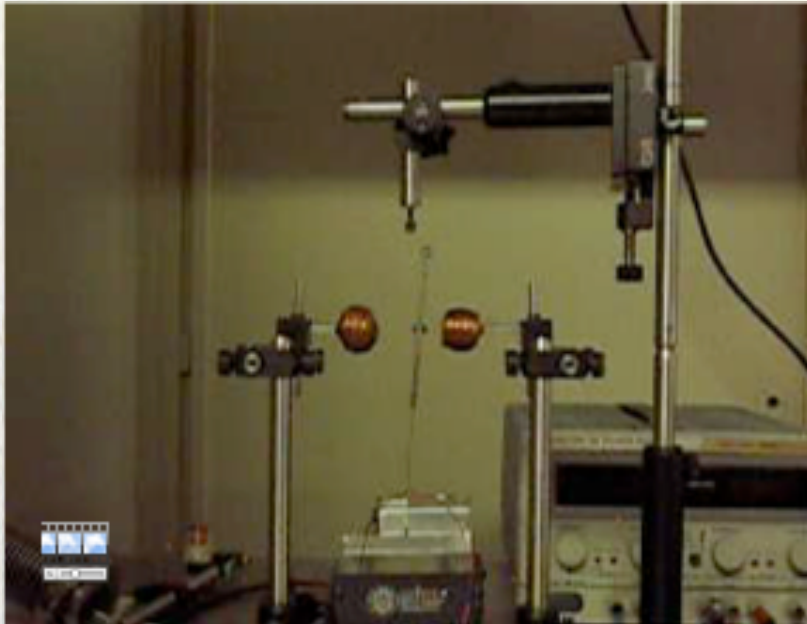
Vibrations energy harvesting

NON-Linear mechanical oscillators



Vibrations energy harvesting

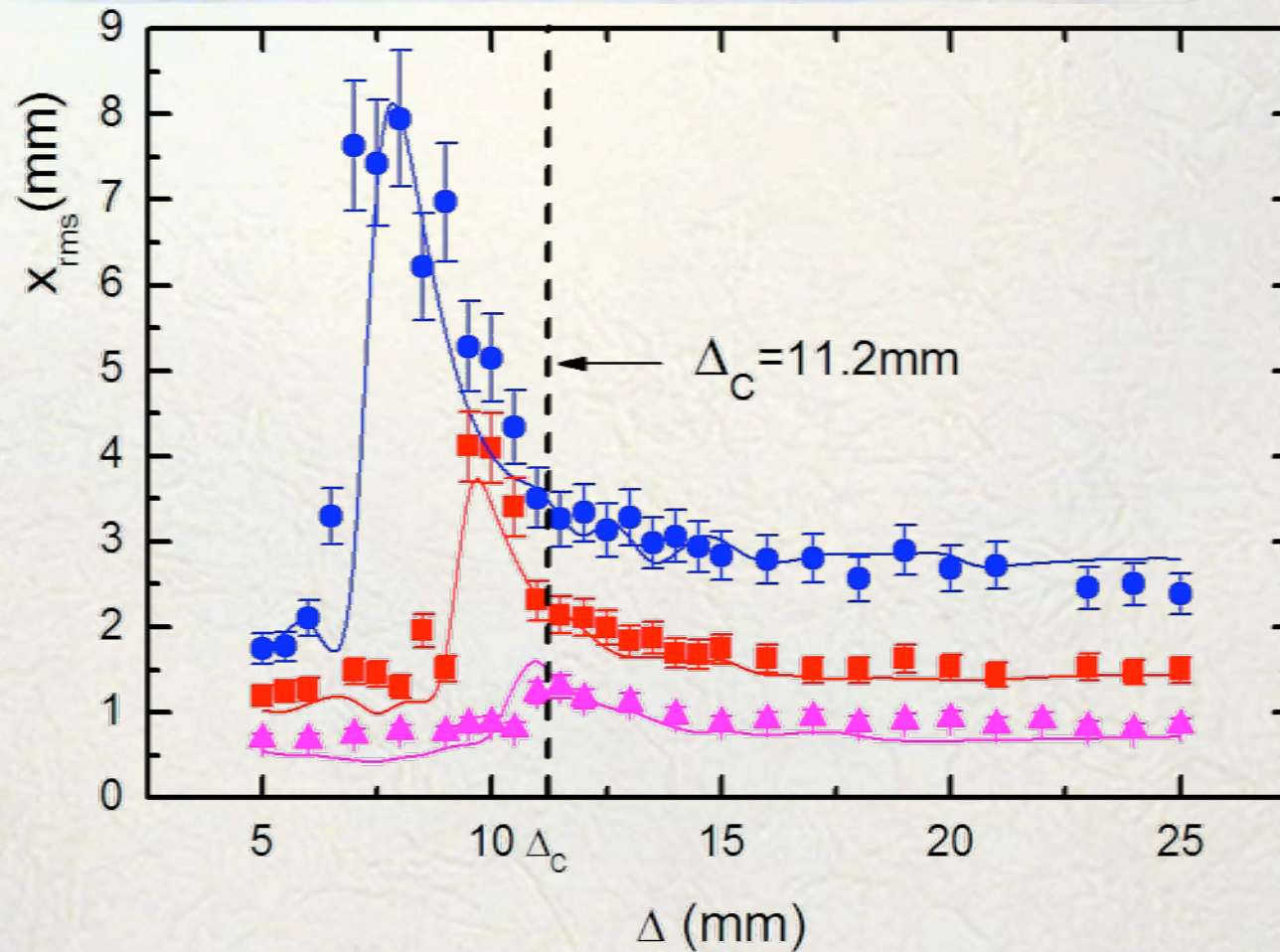
NON-Linear mechanical oscillators



<http://www.nipslab.org/node/1676>

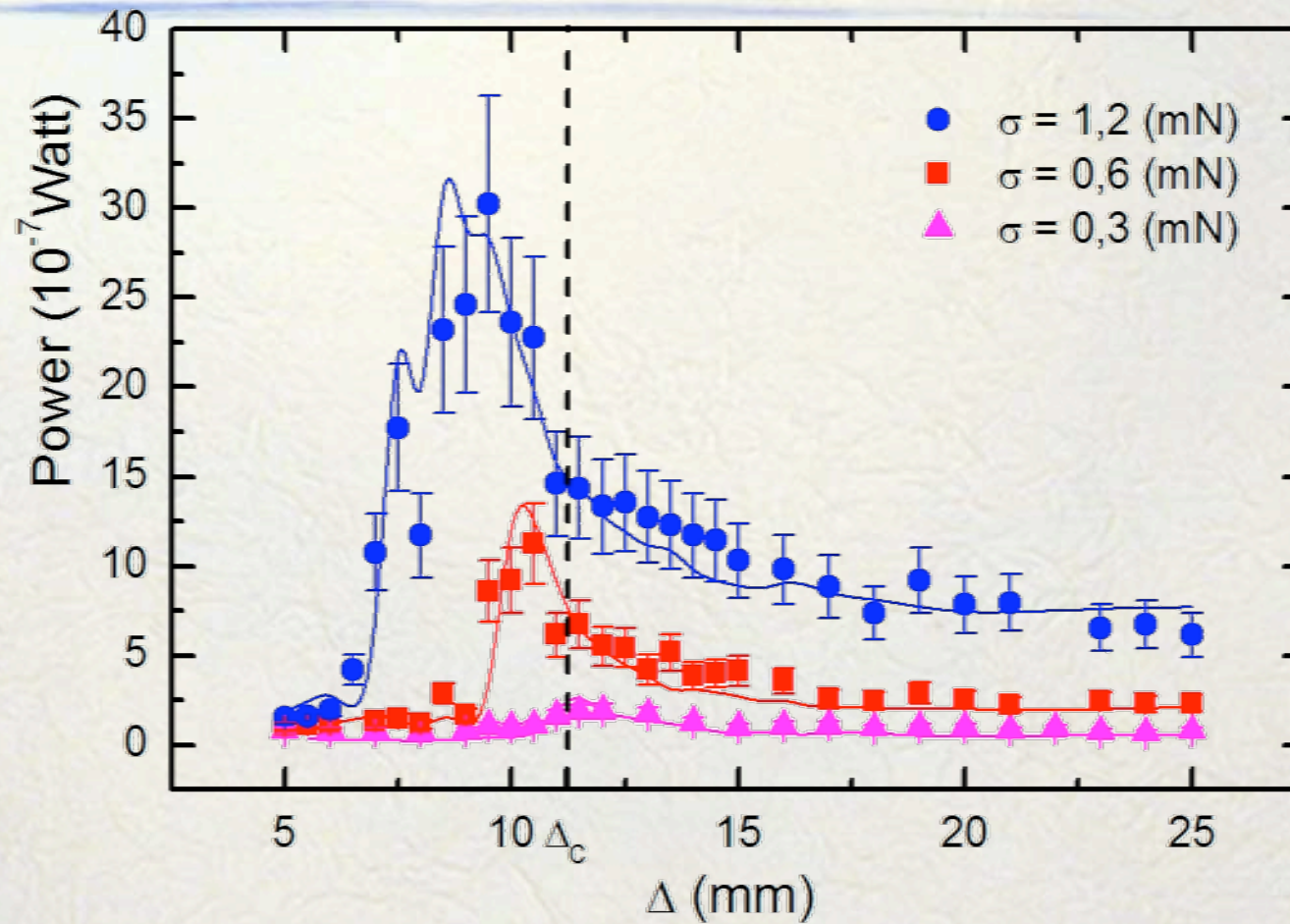
Vibrations energy harvesting

NON-Linear mechanical oscillators



Vibrations energy harvesting

NON-Linear mechanical oscillators



To think about...

- 1) Non resonant (i.e. non-linear) mechanical oscillators can outperform resonant (i.e. linear) ones*
- 2) Non-linear systems are more difficult to treat
- 3) Bistability is not the only nonlinearity available... see:
L. Gammaitoni, I. Neri, H. Vocca, Appl. Phys. Lett. 94, 164102 (2009)